

Prediction of turbulence-cascade interaction noise using modal approach

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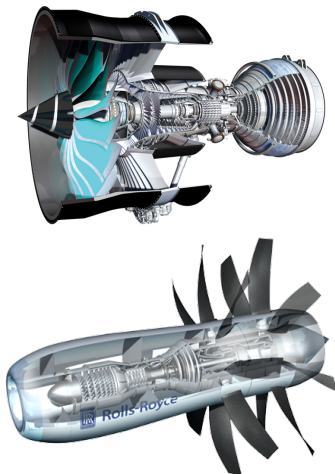
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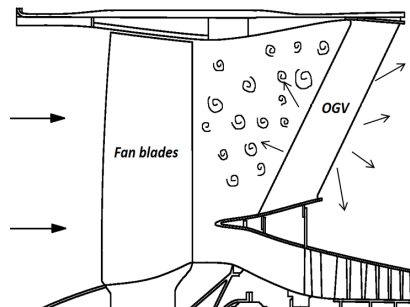
Turbulence-cascade interaction noise

- Fan is one of the dominant noise sources of turbofan engine
- Rotor-OGV interaction is the dominant fan broadband (BB) mechanism
- Ultra high bypass ratio engine
 - Low speed fan
 - More loading on the rotor blades
 - Increase in turbulence
 - Shorter nacelle



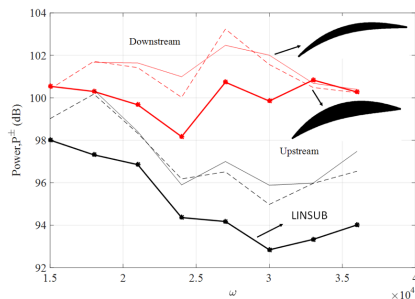
Aim & Objective

- To synthesize rotor wake turbulence as input for linearised Navier-Stokes solver and develop a prediction method for Rotor-OGV interaction noise



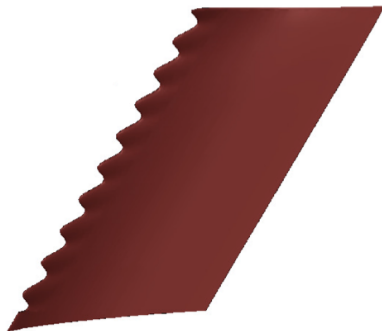
Turbulence synthesis

- Rolls-Royce *p/c*. LNS solver
HYDRA is used
- Frequency domain approach
- Turbulence synthesised by sum of Fourier modes in 2D has been demonstrated
- Geometry effects on cascade interaction noise has been performed



Nature of the flow

- Homogeneous in θ direction
- Inhomogeneous in radial direction
- Anisotropic turbulence
- Varying length scales
- Span-wise correlation effects need to be incorporated
- Correct two-point statistics are important for noise



Need to develop a generalised global approach to establish correct two-point statistics

3D Turbulence synthesis

Fan broadband noise is completely determined by the turbulence velocity cross-spectrum at the leading edge of OGV.

$$C(r, \theta, r', \theta', \omega) = \lim_{T \rightarrow \infty} \left\{ \frac{1}{T} E[u(r, \theta, \omega) u^*(r', \theta', \omega)] \right\}$$

The main objective is to reconstruct the cross-spectral matrix by the modal summation:

$$u(r, \theta, \omega) = \sum_m^{\infty} u_m(r, \omega) e^{im\theta}$$

Each azimuthal modal component $u_m(r, \omega)$ comprises of the sum of radial components n with amplitudes $\hat{u}_{mn}(r, \omega)$, whose radial variation is described by non-dimensional functions (mode shape) $\phi_{mn}(r, \omega)$.

$$u_m(r, \omega) = \sum_n^{\infty} \hat{u}_{mn}(\omega) \phi_{mn}(r, \omega)$$

3D Turbulence synthesis

The cross-spectrum can be expressed in the theta-direction as only a function of $\Delta\theta = \theta - \theta'$, as

$$C(r, r', \Delta\theta, \omega) = \sum_{m=0}^{\infty} C_m(r, r', \omega) e^{im\Delta\theta}$$

$$C_m(r, r', \omega) = \lim_{T \rightarrow \infty} \left\{ \frac{1}{T} E[u_m(r, \omega) u_m^*(r', \omega)] \right\}$$

The cross-spectra C_m in the above series expansion may be obtained by

$$C_m(r, r', \omega) = \frac{1}{2\pi} \int_0^{2\pi} C(r, r', \Delta\theta, \omega) e^{-im\Delta\theta} d\Delta\theta$$

3D Turbulence synthesis

The modal sum is taken over a finite number of radial modes N and then can be expressed in the matrix form

$$\mathbf{u}_m = \Phi_m \hat{\mathbf{u}}_m$$

The cross-spectral matrix C_m , at a number of discrete radial locations for a fixed azimuthal mode number m may be expressed as:

$$\mathbf{C}_m = \Phi_m \hat{\mathbf{C}}_m \Phi_m^H$$

$$\hat{\mathbf{C}}_m = \lim_{T \rightarrow \infty} \left\{ \frac{1}{T} E[\hat{\mathbf{u}}_m \hat{\mathbf{u}}_m^H] \right\}$$

3D Turbulence synthesis ($\mathbf{C}_m = \Phi_m \hat{\mathbf{C}}_m \Phi^H$)

The objective now is to estimate the optimum least-square cross-spectral amplitude matrix $\hat{\mathbf{C}}_m$ for a fixed azimuthal mode m that provides the least squares fit to cross-spectral matrix \mathbf{C}_m .

The optimum cross-spectral $E[\hat{\mathbf{u}}_m \hat{\mathbf{u}}_m^H]$ of vortical mode amplitudes may therefore be written as

$$\hat{\mathbf{C}}_m = \Phi^+ \mathbf{C}_m \Phi^{+H}$$

The upstream and downstream modal acoustic sound power can be expressed in terms of $\hat{\mathbf{C}}_m$ and scattering matrix \mathbf{S}_m^\pm (from LNS solver).

Modelling of Cross-spectral matrix

Sample cross-spectra matrix $C(r, r', \Delta\theta, \omega)$

For simplicity we first assume the cross-spectrum follows Gaussian distribution and turbulence is assumed to be locally homogeneous.

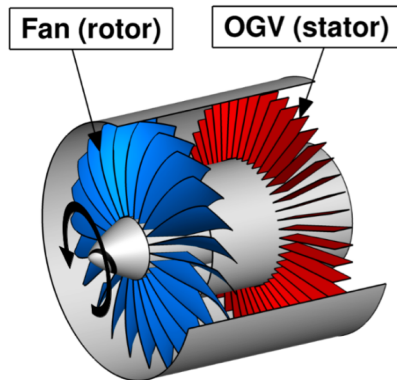
$$C(r, r', \Delta\theta, \omega) = \left(S(r, \omega) S(r', \omega) \right)^{0.5} \gamma(\Lambda_\theta, \Lambda_r, r - r', r(\Delta\theta))$$

$$\gamma(\Lambda_\theta, \Lambda_r, r - r', r(\Delta\theta)) = e^{-(r-r')/\Lambda_r^2(\omega)} e^{-r(\Delta\theta)/\Lambda_\theta^2(\omega)}$$

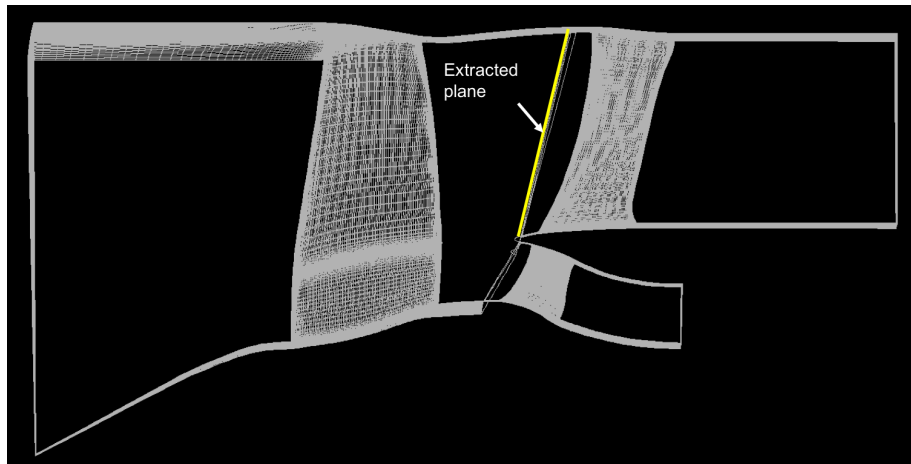
Length scales and intensity variation along the radius are extracted from
RANS

RANS mean flow

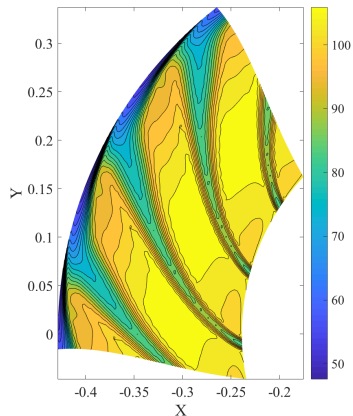
- Mesh size: ≈ 6.5 million cells
- Fully resolved boundary layer ($y^+ < 1$) on blades
- RANS with mixing-plane between rotor and stator blocks
- k- ω turbulence model used
- AP: AP geometry w/tip clearance of 0.78 mm
- CB, SL: hot geometry w/tip clearance of 0.2 mm



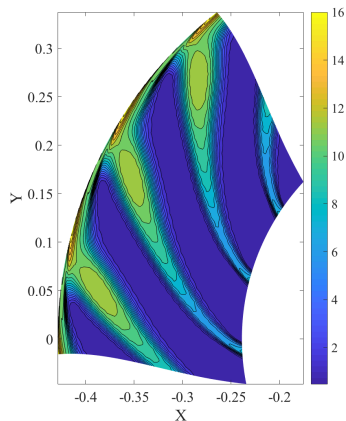
Extraction plane



Axial velocity and turbulent intensity contours

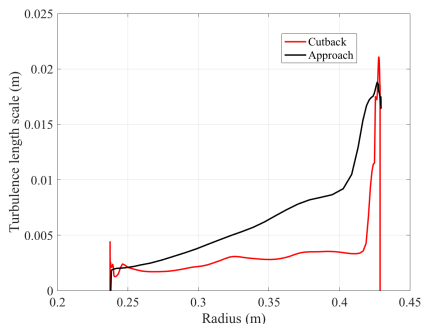


Axial velocity (m/s)

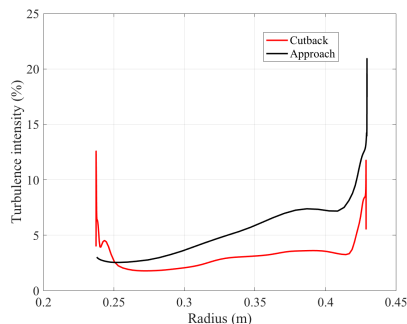


Turbulence intensity (%)

Circumferentially averaged turbulence intensity and length scales



Turbulent length scale



Turbulence intensity

Vortical mode shapes

Choice of vortical mode shapes

Our work has so far considered a total three vortical modal solutions.

- Normal modes of LEE, Φ_{LEE}
- Fourier modes, Φ_F
- POD modes, Φ_{POD}

Important requirements:

- Divergence free
- Efficient representation of cross-spectral matrix

Normal modes of LEE (Φ_{LEE})

Linearized Euler equations can be expressed in matrix form as

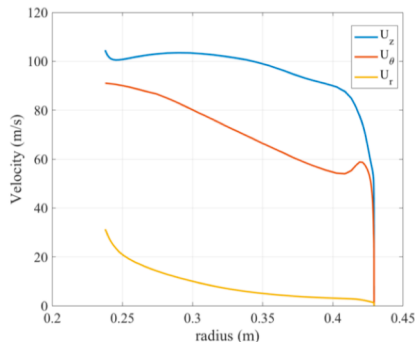
$$A \frac{\partial}{\partial z} \begin{bmatrix} u_r \\ u_\theta \\ u_z \\ p \end{bmatrix} = B \begin{bmatrix} u_r \\ u_\theta \\ u_z \\ p \end{bmatrix}$$

$$A = \begin{bmatrix} U_z & 0 & 0 & 0 \\ 0 & U_z & 0 & 0 \\ 0 & 0 & U_z & 1 \\ 0 & 0 & 1 & U_z \end{bmatrix}$$

$$B = \begin{bmatrix} -i\omega + \frac{imU_a}{r} & \frac{U_a}{r} & 0 & -D_r \\ -\frac{U_a}{r} - (D_r U_\theta) & -i\omega + \frac{imU_a}{r} & 0 & \frac{im}{r} \\ -(D_r U_z) & 0 & -i\omega + \frac{imU_a}{r} & 0 \\ -\frac{1}{r} - D_r & \frac{im}{r} & 0 & -i\omega + \frac{imU_a}{r} \end{bmatrix}$$

where

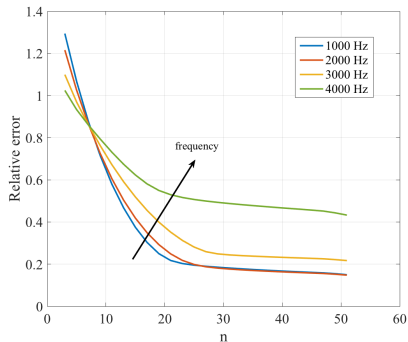
Solution to the Eigenvalue equation are the normal modes of LEE (Φ_{LEE})



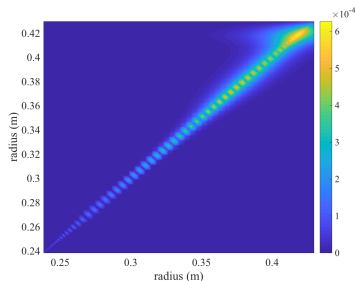
Fourier modes (Φ_F)

$$\phi_{F,mn}(r, \omega) = e^{ik_r r} \rightarrow \hat{\mathbf{C}}_m = \mathbf{\Phi}^+ \mathbf{C}_m \mathbf{\Phi}^H$$

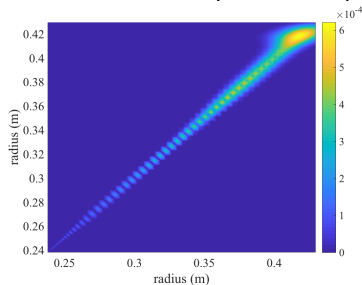
$$\text{error} = \frac{\|C_m - C_{m,\text{new}}\|}{\|C_{m,\text{new}}\|}$$



Original



Reconstructed (20 modes)

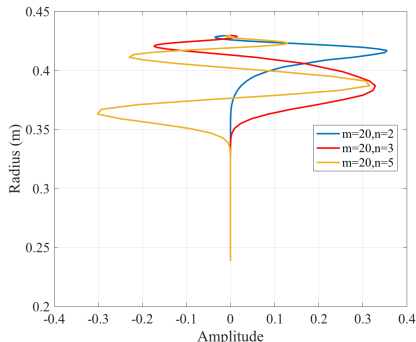


POD modes (Φ_{POD})

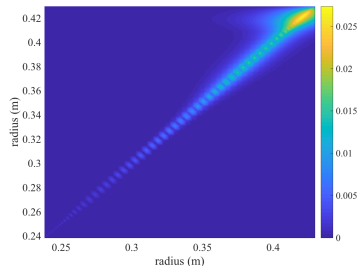
The POD of the two-point velocity cross-spectrum is of the form,

$$C_m(r, r', \omega) = \sum_{n=1}^N \lambda_{mn}(\omega) \phi_{mn}(r, \omega) \phi_{mn}^*(r', \omega)$$

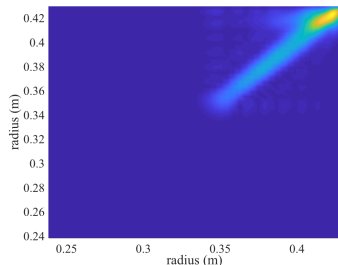
$$\hat{\mathbf{C}}_m = \text{diag}(\lambda_{m1}, \lambda_{m2}, \lambda_{m3}, \dots, \lambda_{mn})$$



Original



Reconstructed (8 modes)



Comparison

Approach	Advantages ☺	Disadvantages ☹
Φ_{LEE}	Naturally divergence-free	Sharp modes Badly conditioned
$\Phi_{Fourier}$	Good reconstruction of wake Easy to impose divergence-free	More radial modes are required Expensive to capture large scale structures
Φ_{POD}	Good reconstruction of large scale structures with limited modes	Expensive for wake turbulence Few modes are badly discretized

Different combinations can be used for different operating conditions

Conclusions

- Presented a generalised framework for globally reconstructing the cross-spectral matrix
- Three different vortical modes are investigated to reconstruct the exact cross-spectral matrix: 1) Fourier modes, 2) POD modes 3) Normal modes of LEE.

Future work

- Development of time domain approach as frequency domain is expensive for Broadband noise calculations.

Thank you!

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