

Modern numerical and analytical methods in duct acoustic simulations

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Can CAA really be used in a numerical experiment to replace real experiments today?

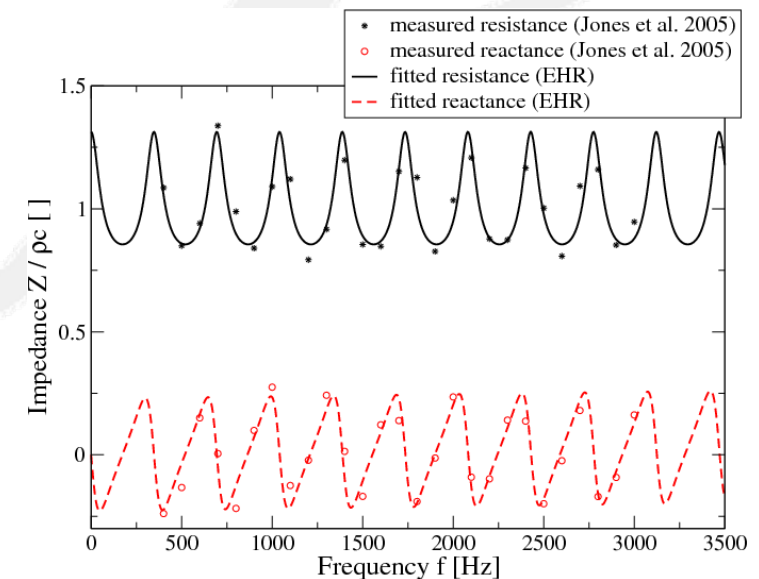
- What is possible with time-domain CAA, today? – An inventory from the developers:
 - Complex geometries can be handled with some effort
 - E.g. using Chimera approaches in connection with Finite Difference methods
 - Very detailed approximations of the base flow conditions (incl. boundary layers, shear layers and choked flows) can be handled by CAA.
 - Development of ideally non-reflective boundary conditions is not yet finished – often reflections are smaller than with “anechoic terminations” in the experiment.
 - Noise reduction devices (acoustic lining and connected systems of resonators) can be modeled in the time domain
 - There are still many questions in this area
 - An accurate link between model and liner hardware is provided only by experiments
 - Main advantage of the time domain simulation: Arbitrary source signals are possible without extra effort. – This is rarely studied – Why?
- High-fidelity simulations for all ducts of an aeroengine are possible! –
Can this be used for numerical experiments?
 - Requires methods for analysis and interpretation of the CAA result.
 - What is then required to make mechanized targeted optimization with CAA methods possible.

- Time domain CAA methods and related issues
 - Time domain impedance boundary conditions
 - Handling of complex geometries
 - Results validation and reliability
 - Tools for interpreting the results
 - Optimization and simulation time issues
- Selected results are presented inline with the related methods

Time domain impedance boundary conditions

Time-domain impedance boundary conditions (TDIBC's)

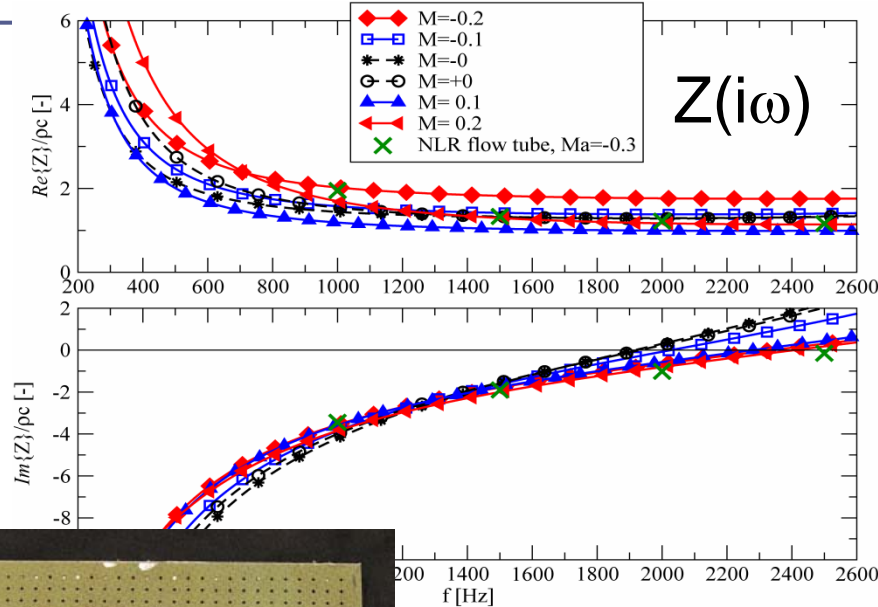
- TDIBC's are a generalization of the hard wall and fully non-reflective boundary conditions.
- They are applicable with CAA as well as for CFD
- Can model acoustic lining, soft walls, connected resonators or the frequency response of connected ducts instead of anechoic properties!
- Preconditions for the application of TDIBC's:
 - How to obtain the model parameters for a specific liner?
 - Model or resolve boundary layer effect on the impedance of the lined wall?
 - Time-domain impedance boundary conditions require a model for the frequency response of a liner – Can all types of liner be covered?



The Extended Helmholtz Resonator [Rienstra, 2006] is applied here to model the frequency response of a liner → Only single Helmholtz Resonator, but can fit termination impedance of NASA flow duct over large frequency range!

A way to obtain TDIBC-model parameters – Impedance Education from measurements

Model parameters and $Z(i\omega) = ?$



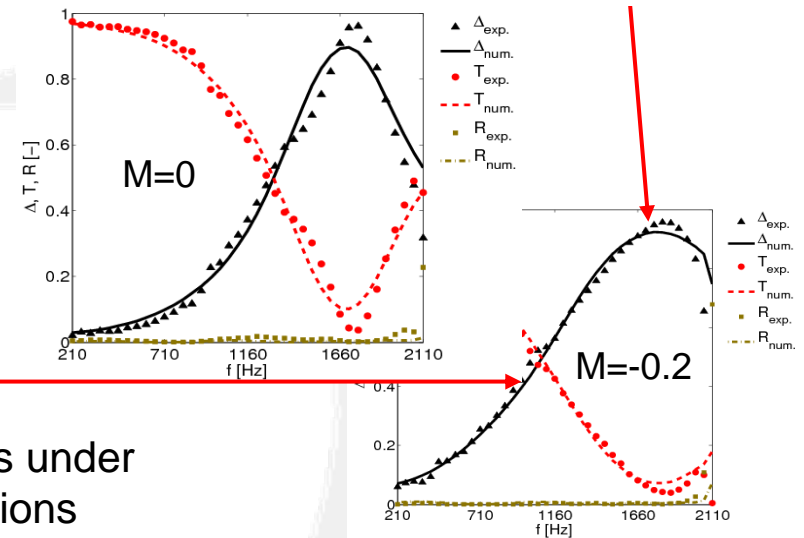
Model parameters

Optimization (MATLAB)

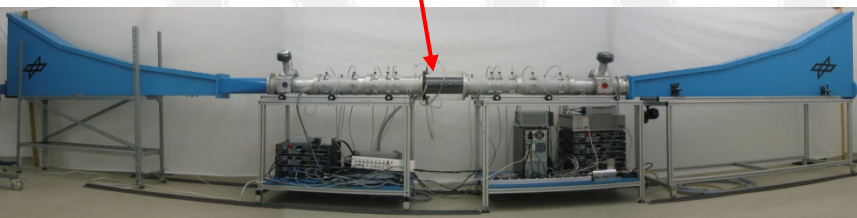
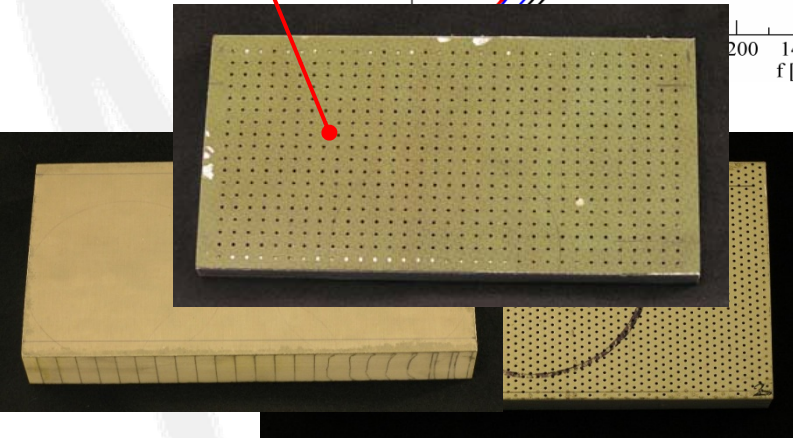
MATLAB Optimization tool calls TUBA, Variation of the model parameters, Objective: Fit measurement result

Result: Parameters and Impedance of a liner under relevant flow conditions

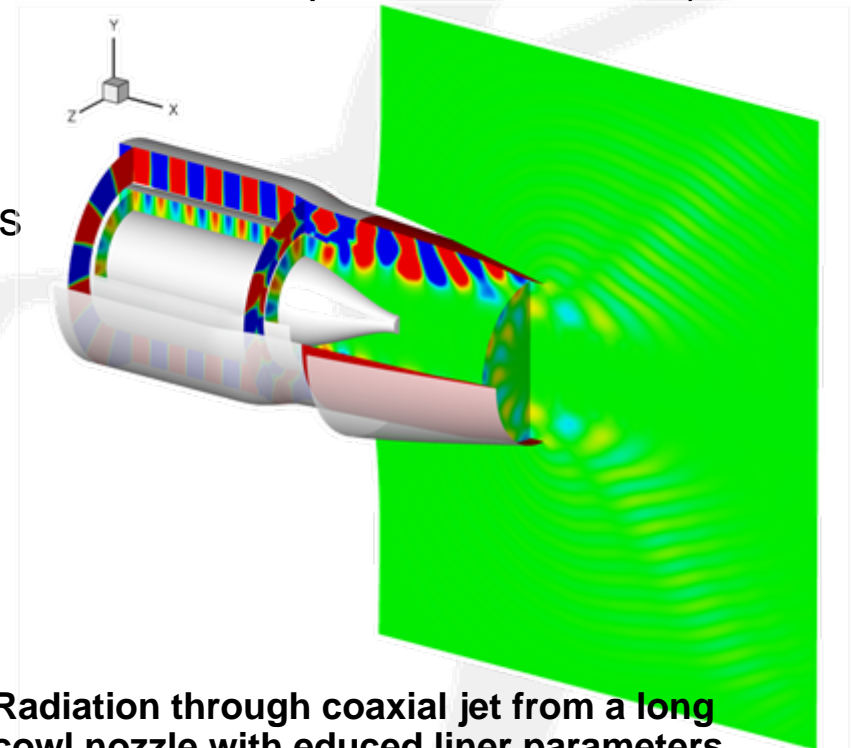
Multifrequency excitation, Grid size <20000 cells, Turn around time ~5...90 min



Measurement of reflection and transmission factors under relevant flow conditions



- The **impedance eduction** in the time domain provides a set of parameters for the impedance model
 - This parameter-set is best approximating the experimental observation
 - Parameters describe the frequency response of the impedance.
 - Parameters provide a relation between physical hardware (e.g. geometry) and measurement result under flow conditions
(eventually, only if a physical interpretation of these parameters exists)
- The result allows a predictive simulation of the liner effect under application conditions:
 - Presuming that the measurement can reproduce all necessary basic conditions (flow velocity, temperature, frequency) as the practical application
- The result allows to give directions for an improvement of the liner towards the designated impedance

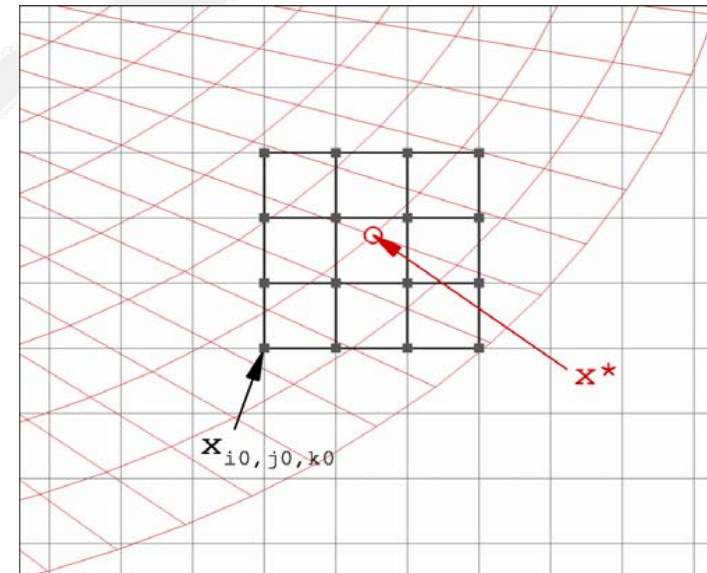


Radiation through coaxial jet from a long cowl nozzle with educed liner parameters

Handling of complex geometries: Chimera technique

- Requirements of standard finite difference block structured CAA methods clearly conflict with the aim of handling complex geometries:
 - Extremely **challenging** and time consuming grid generation in 3D.
 - **Distorted** cells in the mesh have to be avoided, to
 - Keep the accuracy of the method high,
 - Avoid instabilities, and
 - Keep the simulation time acceptable.

- An **overset-grid-method** is applied (Chimera) to avoid penalties due to the meshing of complex geometries:
 - Allows Orthogonal and equidistant host mesh
 - Separate body fitted overset meshes
 - Requires Interpolation between the meshes
 - High-order FD-schemes require **high-order interpolation**, here based on Lagrangian-polynomials on 4x4x4 (5x5x5) points



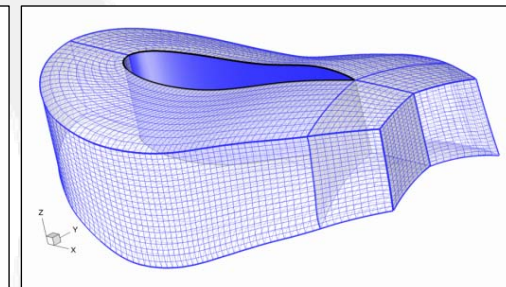
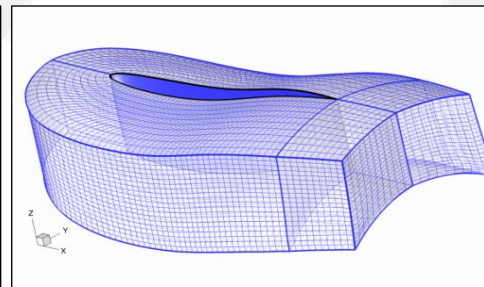
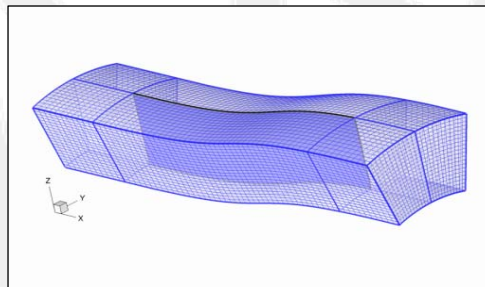
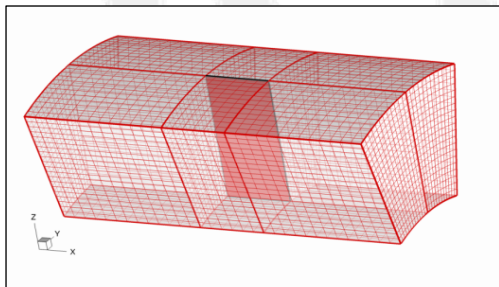
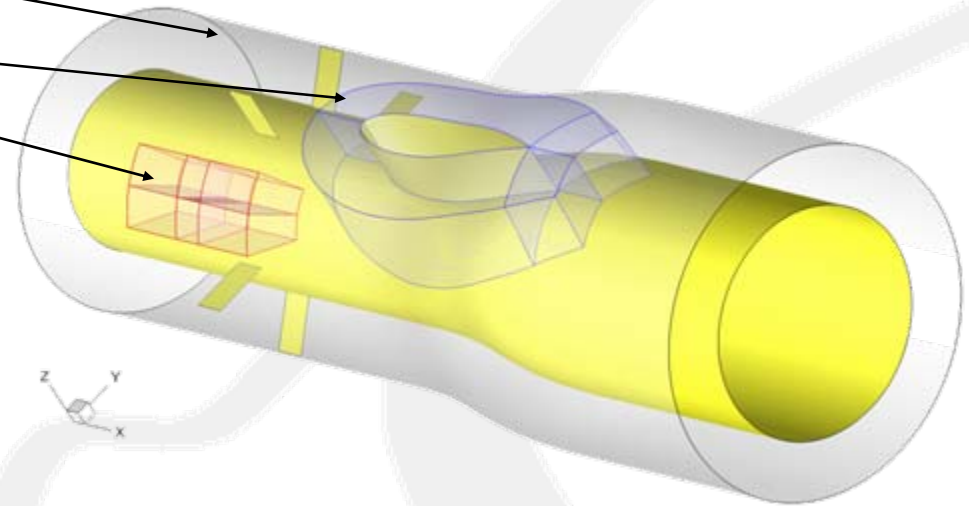
Example of a 2D interpolation stencil

Meshing of an S-duct with installations

➤ Meshing of an S-duct with installations using the Chimera technique

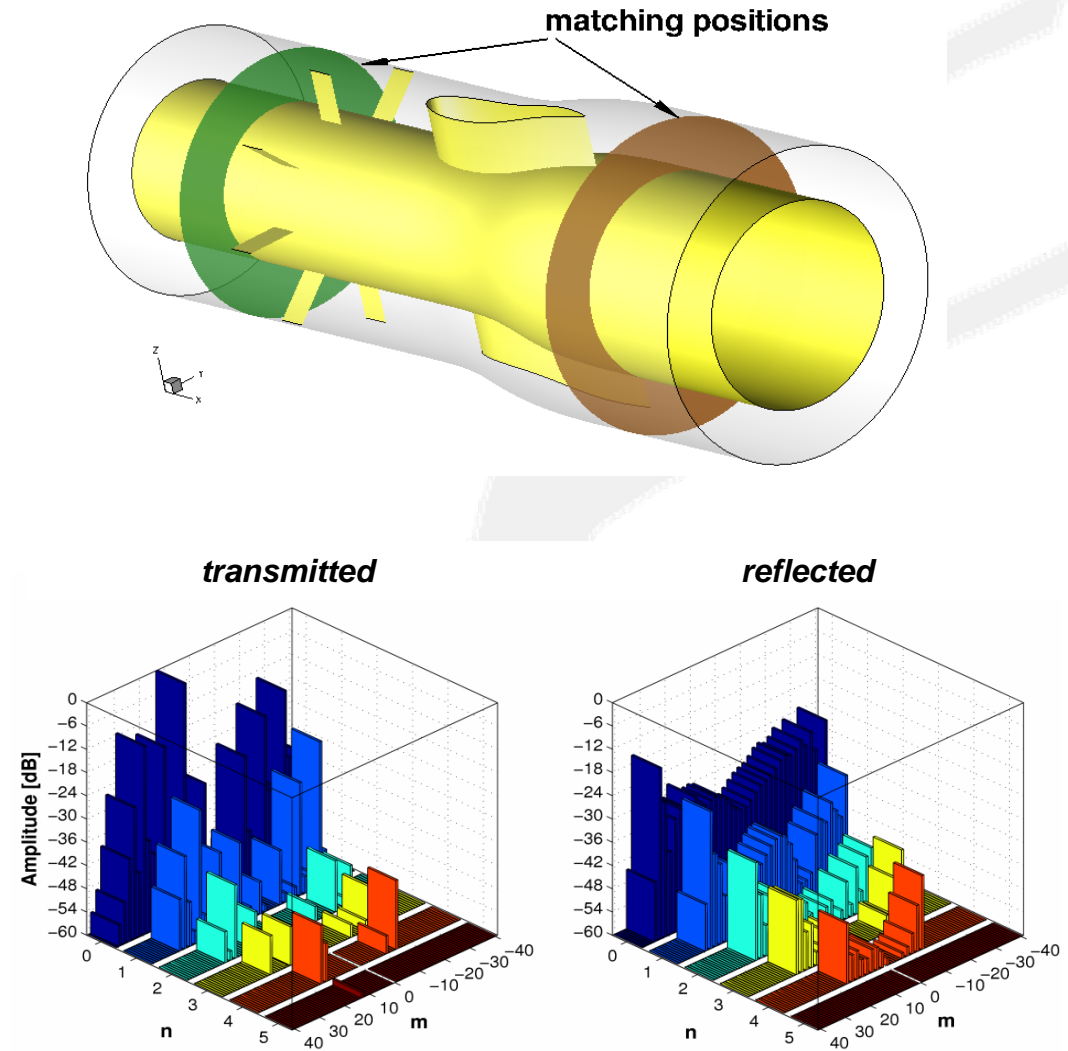
- Host grid:
 - uniform and local orthogonal O-type grid
- Overset grids:
 - Body fitted C-type for profile bifurcations
 - 6-block H-type grid for struts and planar plate bifurcation

- 2.1 – 2.4 Million nodes
- 4 – 12% interpolated nodes
(w and w/o struts)



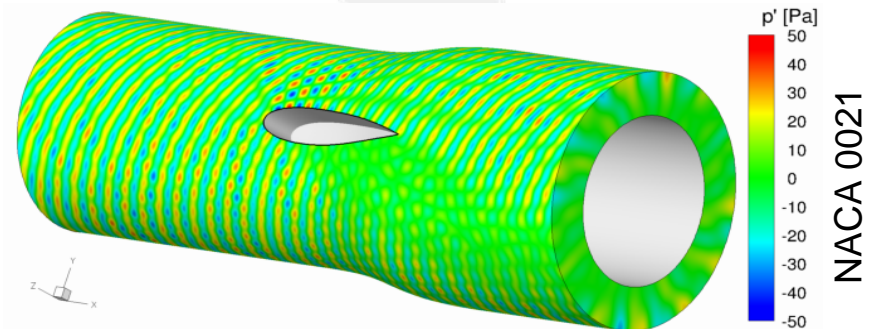
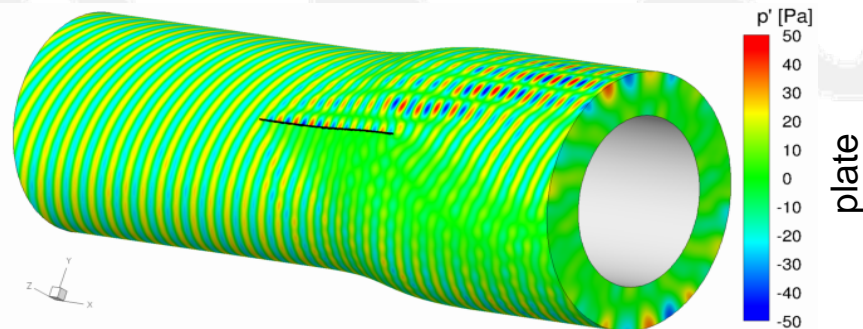
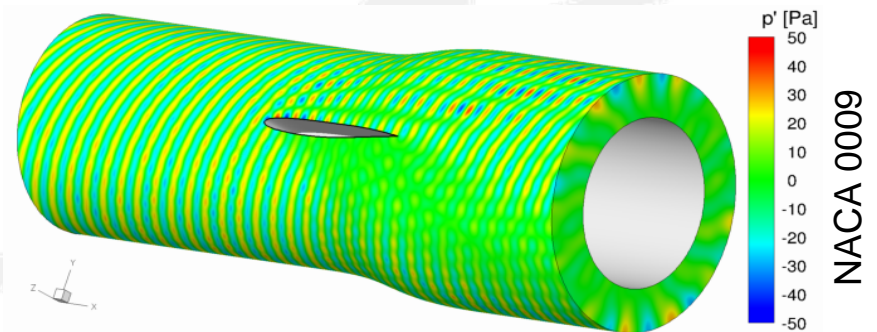
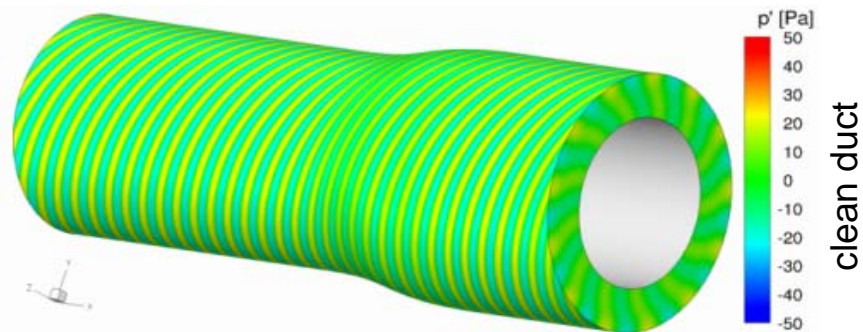
Interpretation of the CAA results

- Analysis with mode-matching method [OVENDEN & RIENSTRA, 2004]
- Originally developed for extraction of acoustic components from unsteady CFD data
- Also suitable for modal analysis of the propagated sound field at arbitrary planar cross-sections in the duct
- Provides amplitudes and phases of all present modes for a given frequency
 - Azimuthal mode m
 - Radial mode n
 - Propagation direction (transmitted, reflected)

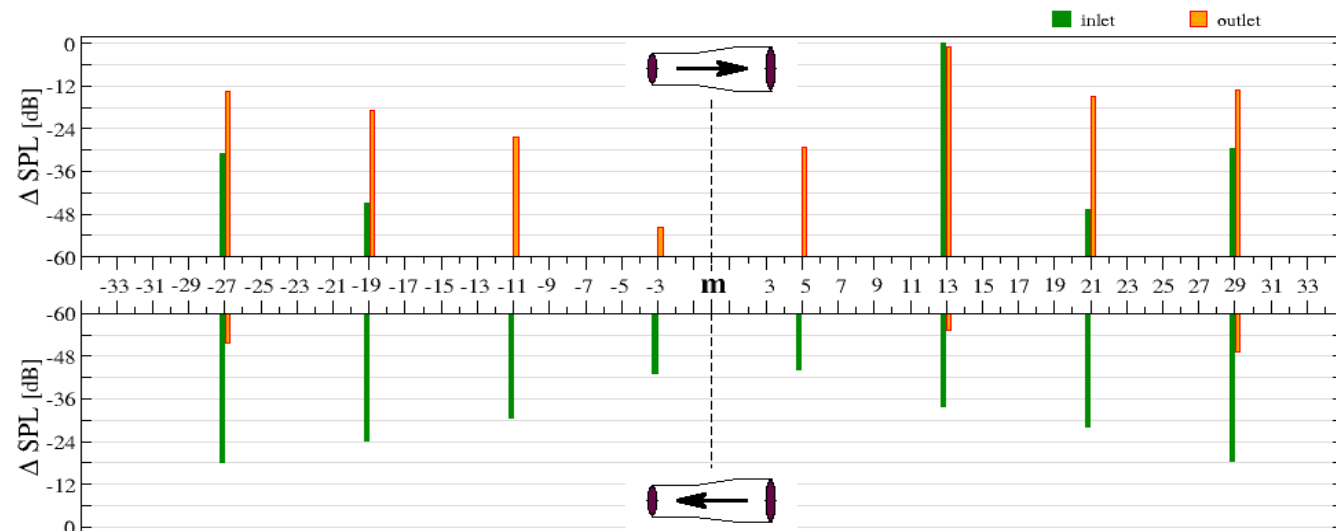
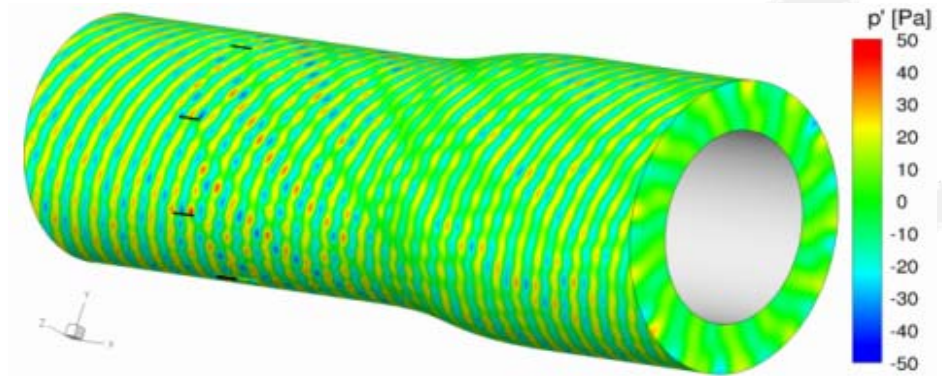


Installation effects in a S-duct – Effect of the profile width

- Splitters disturb the mode propagation
 - Reflections on the exposed side and shadowing on the rear side in propagation direction
- Increasing of the profile width
 - Decreased downstream transmission
 - Increased upstream reflection



- Due to the 8 struts every 8th azimuthal mode is excited: theory of TYLER & SOFRIN (1962)
- Source mode $m=13$ keeps dominant
- Amplitude loss of source mode
 - ~1 dB
 - Cause: smaller number of scattered modes
- Scattered modes at least 12 dB quieter
- Smaller influence compared to splitters

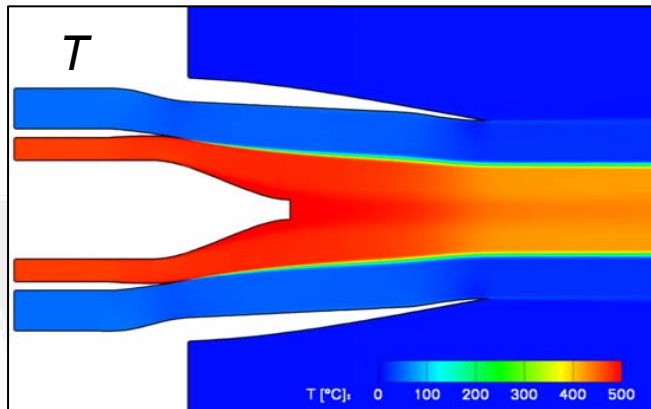


Numerical Experiment: Radiation from a lined coaxial nozzle

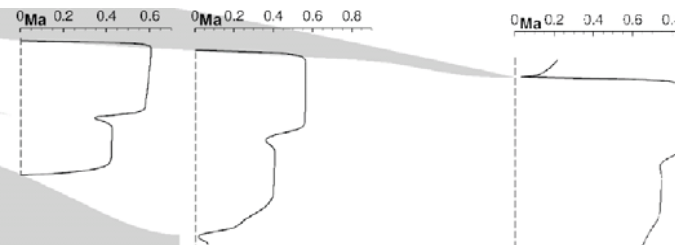
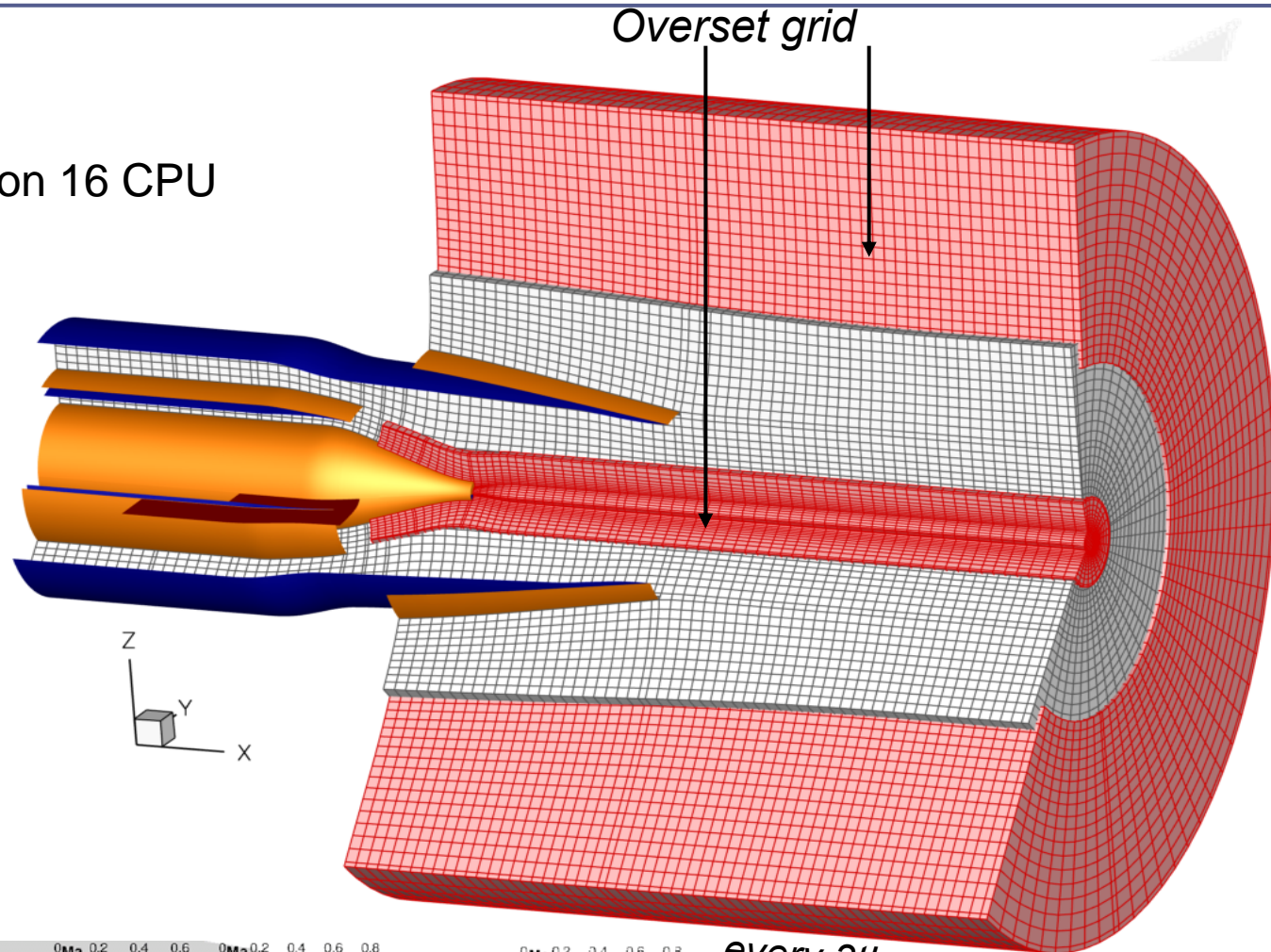
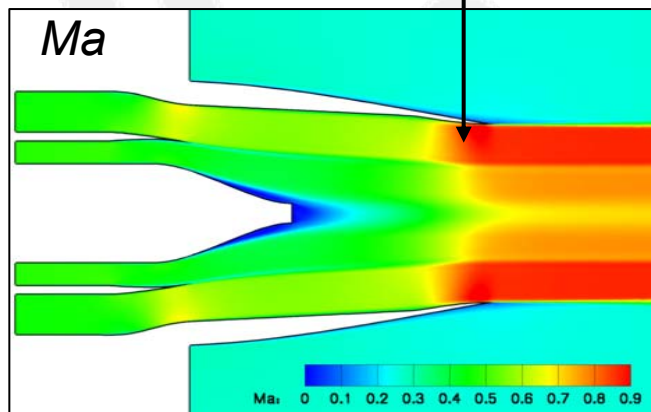
Numerical Experiment: Radiation from lined coaxial nozzle with splitter

➤ Chimera mesh

- $6 \cdot 10^6$ grid points
- 12 h computational time on 16 CPU cores



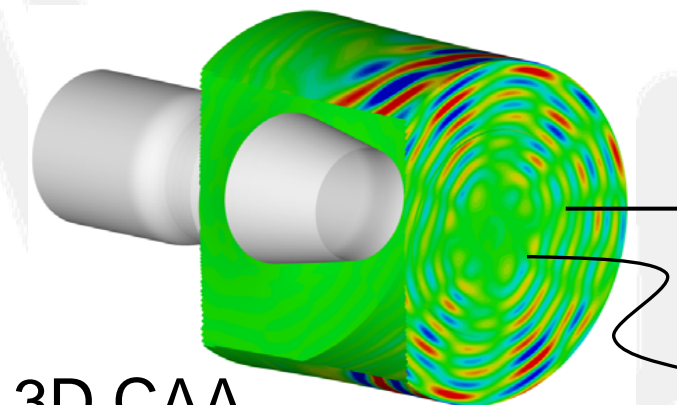
$M_{\max} = 0.9$



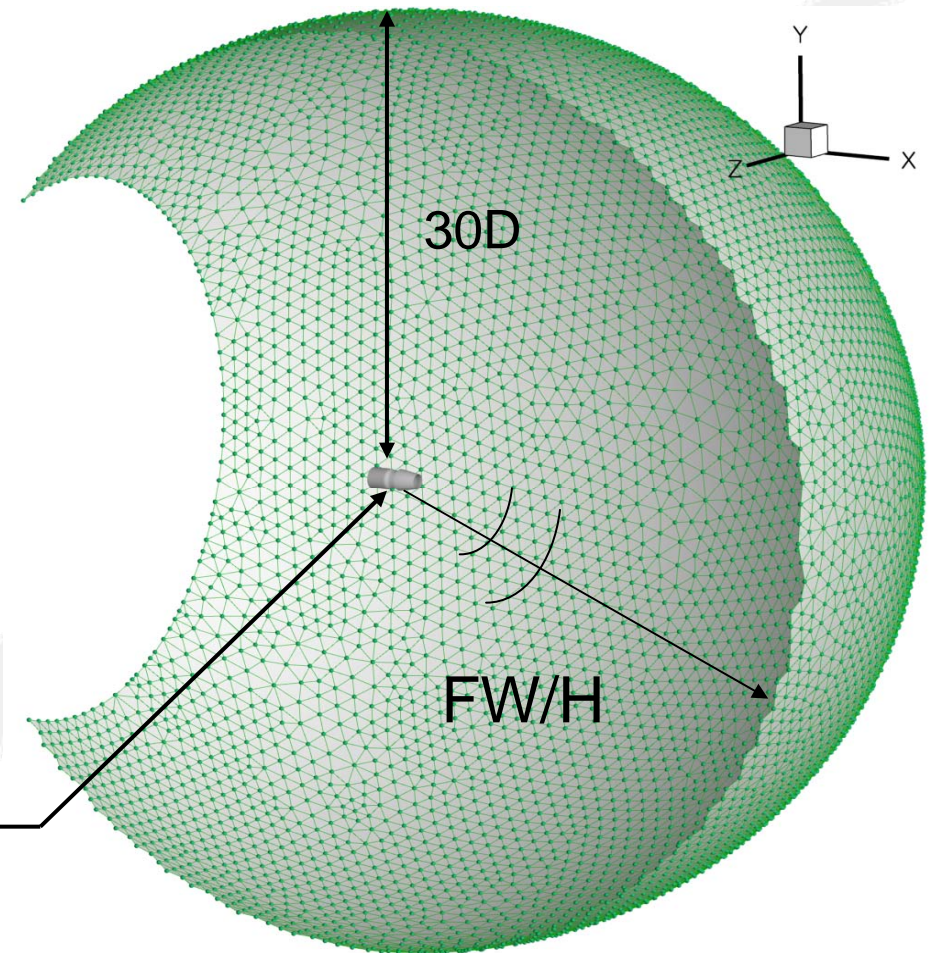
every 3th gridline

Acoustic analogy of Ffowcs-Williams & Hawkings for far field radiation

- Far-field calculation based on spatial and time resolved CAA simulation still impossible with current computing resources.
- Acoustic analogy of Ffowcs-Williams & Hawkings much faster
 - Parallel implementation (MPI)
 - Computational time: **1h** on 4 CPU



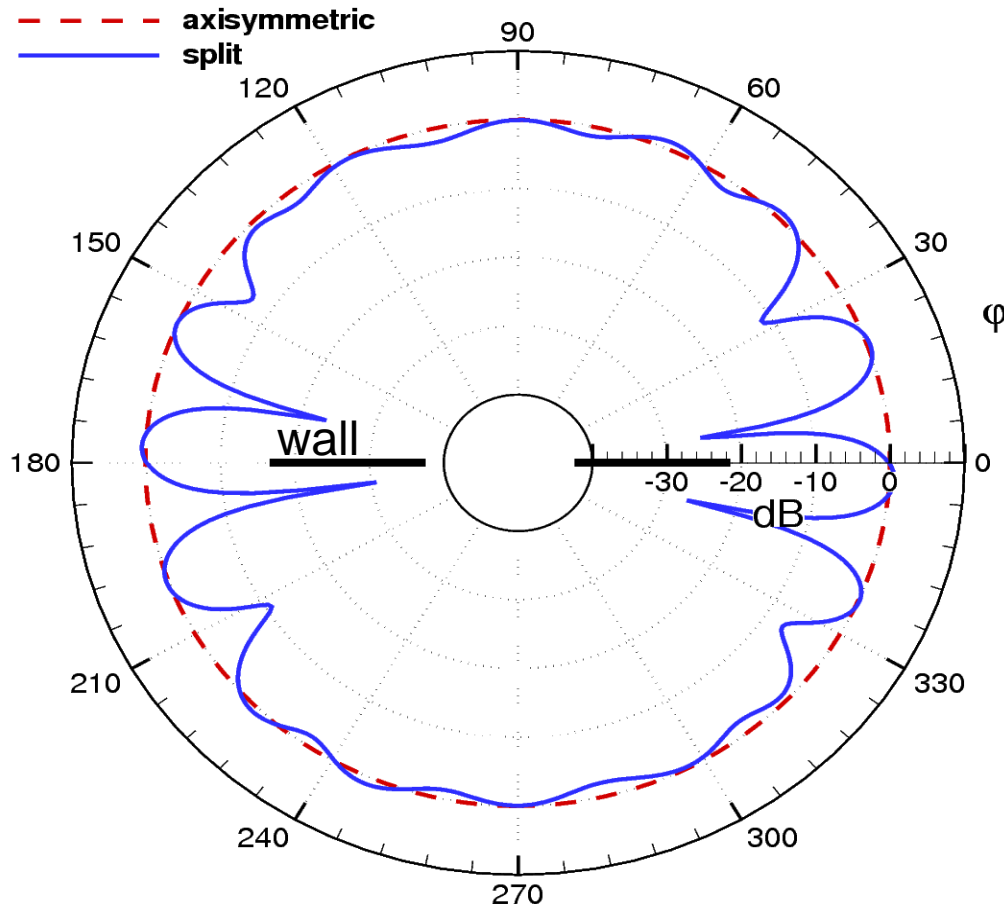
Exchange surface around coaxial exit nozzle ($R = 30 D$)



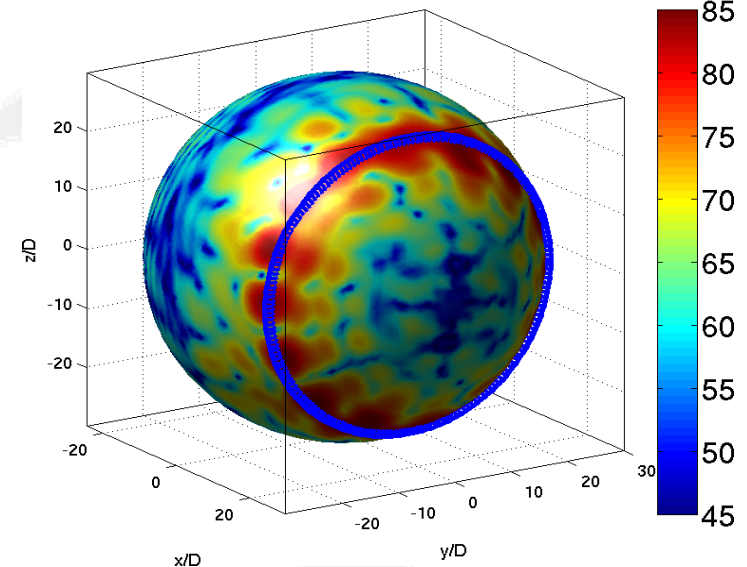
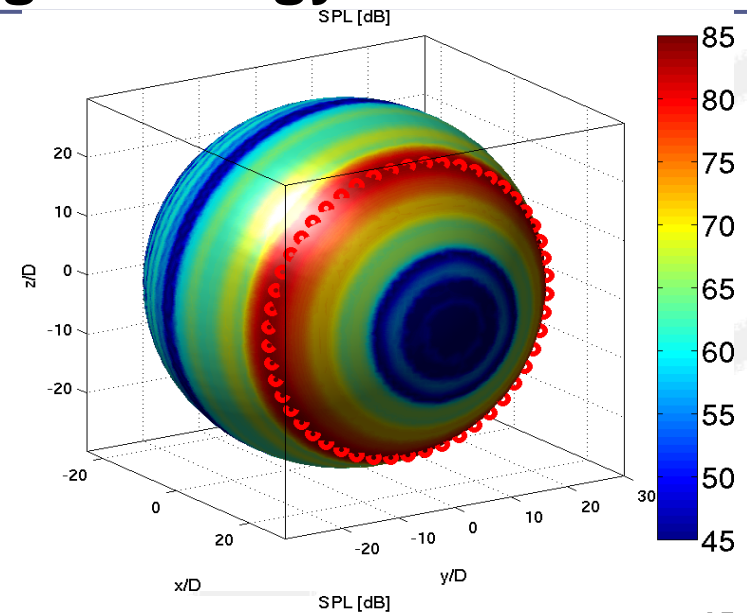
10^4 Observer on sperical surface

Far field radiation characteristics obtained with Ffowcs-Williams & Hawkings Analogy

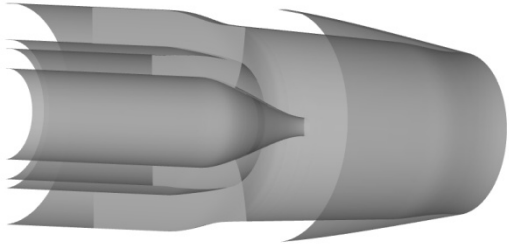
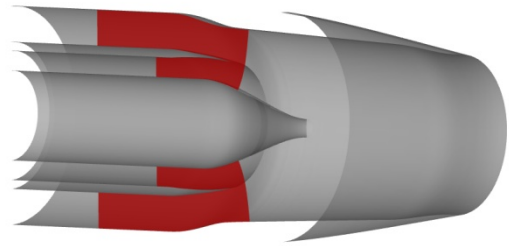
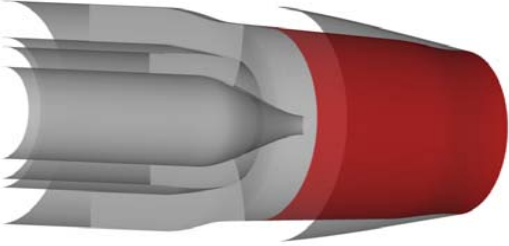
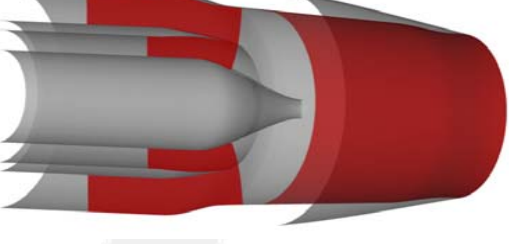
- Bifurcation produces irregular radiation characteristics



Polar angle $\psi=55^\circ$

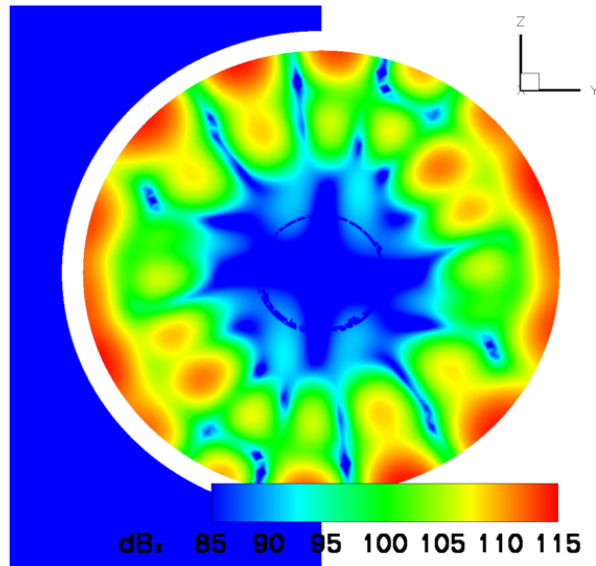
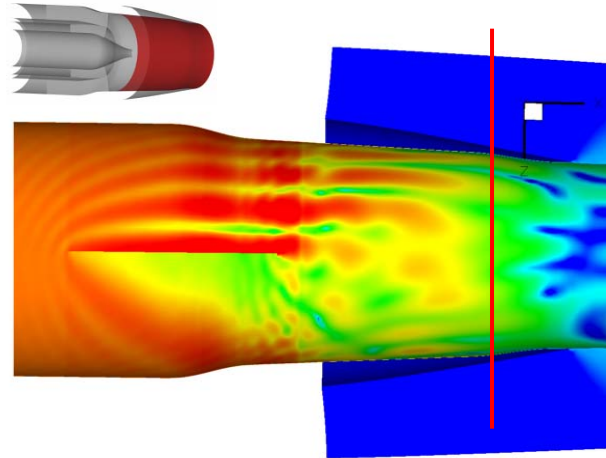


- Extended-Helmholtz-Resonator (**EHR**) Model with 5 Parameters [RIENSTRA, 2006]
- Resolved wall boundary layers and shear layers in the base flow
- **Model parameters are obtained from “impedance eduction”** with an inlet liner from AleniaAermacchi [BUSSE ET AL., 2008],
- Inlet liner not optimized and practical applicable in the bypass duct → numerical experiment!

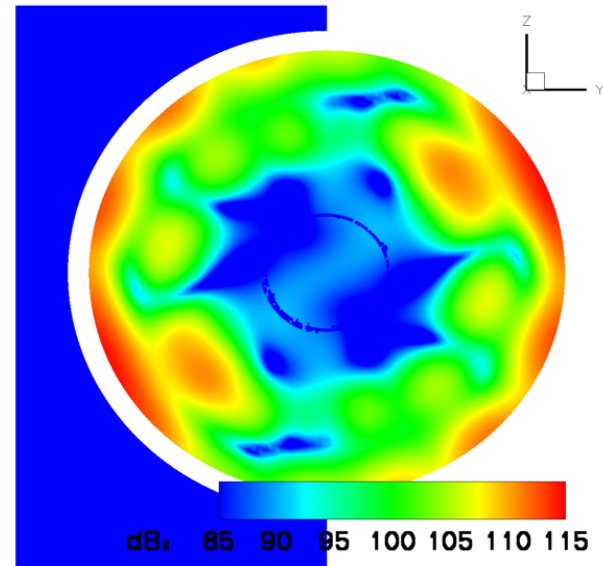
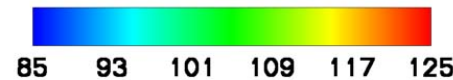
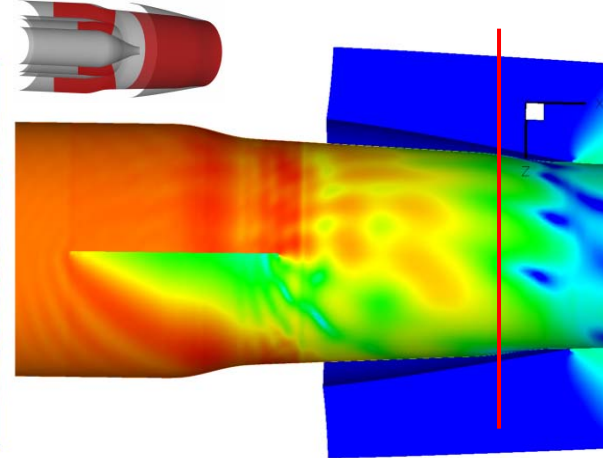
No lining, hard bifurcations (reference)	
Lined bifurcations	
Lined exit nozzle	
Lined exit nozzle and bifurcations	

Application of educed model for inlet liner – Numerical experiment: Radiation from a coaxial nozzle

Only nozzle lined



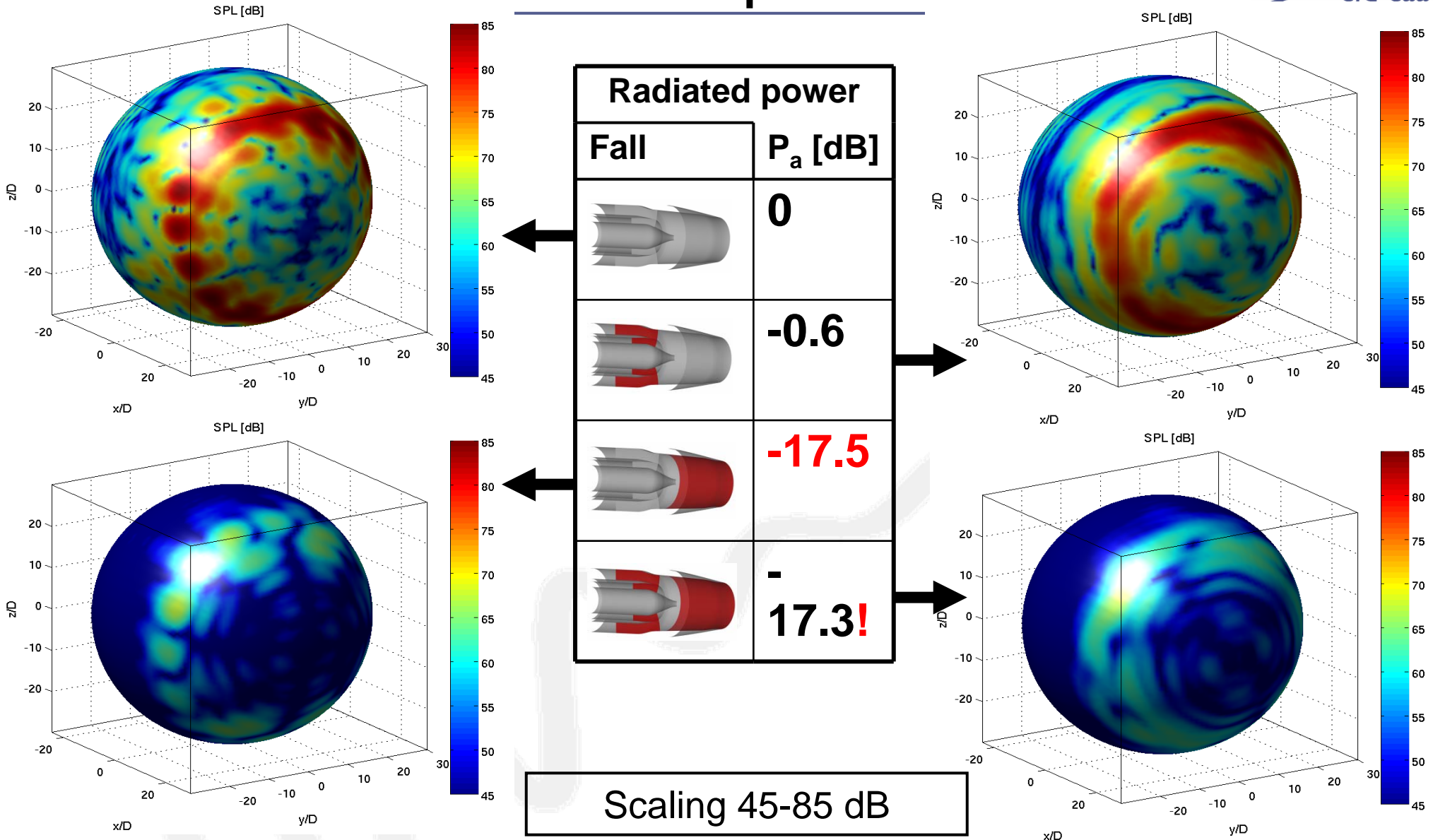
Cut at
 $x = -0.5$



Nozzle and splitter walls
lined with typical inlet liner

P_{RMS} [dB]

Radiation from a lined coaxial nozzle – Radiated power



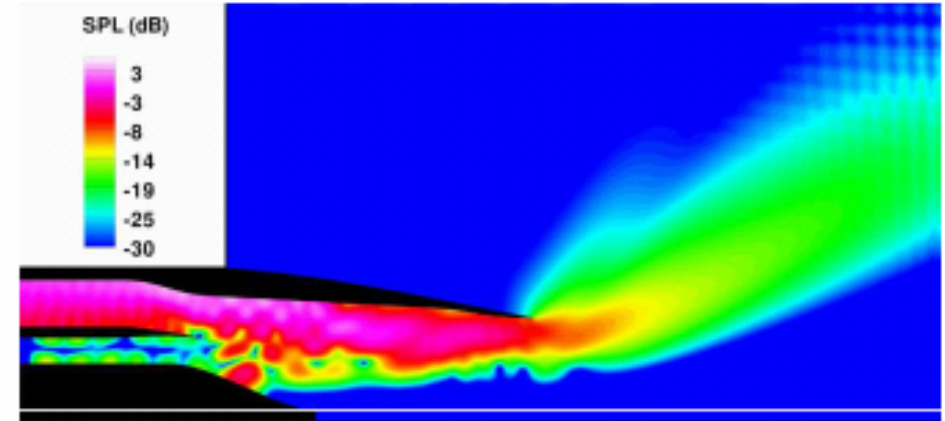
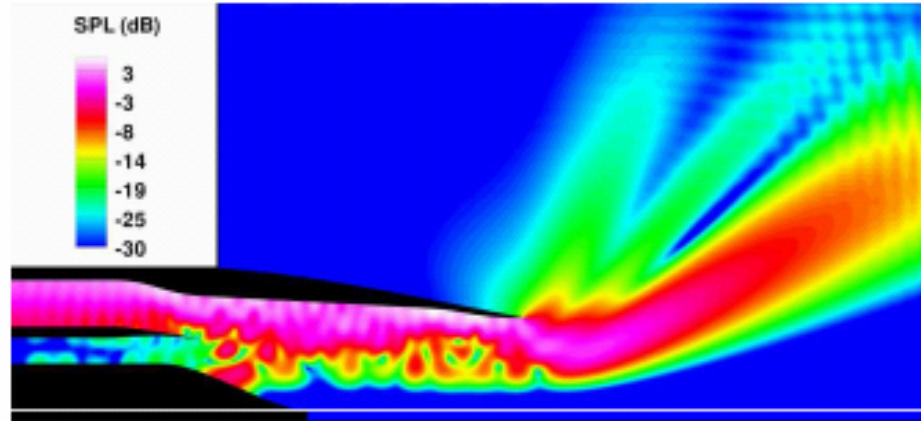
GPU based acceleration of the simulation for optimization purposes

GPU based simulation of the lined coaxial nozzle configuration in 2D

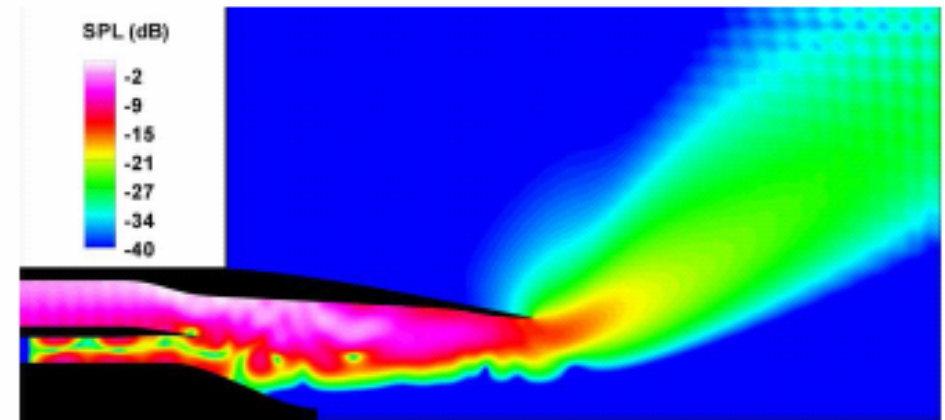
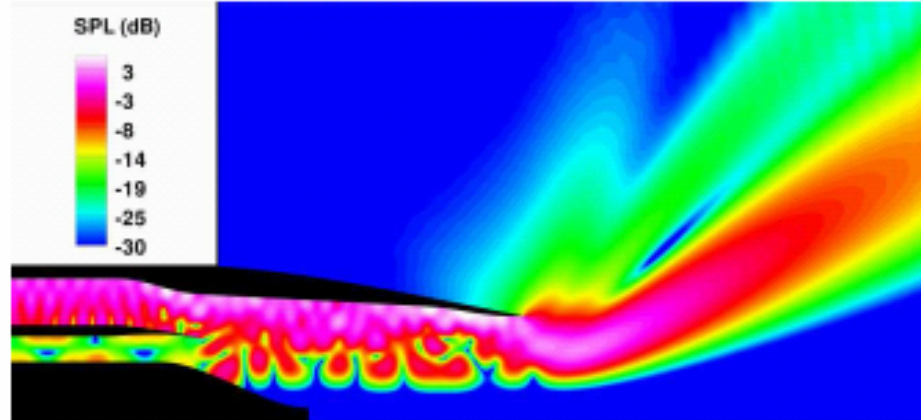
hard-walled

lined

GPU



CPU



➤ CPU and GPU in reasonably agreement (non-reflective BC missing for GPU)

- TUBA3D: 8.560 minutes
- GPU: 22 minutes

→ Speed up of over 300

Large speed-up by GPU computing possible

CPU

GPU

High fidelity University research



Calculable for single configuration in industrial application

Calculable for single configuration in industrial application



Optimizable (calculates multiple configurations)

Optimizable (calculates multiple configurations)



Real time simulation?

GPU computing makes optimization available for many applications at low hardware costs!

Observed speed up requires writing a new code!

Conclusions

- Numerical experiments predicting sound propagation and radiation are possible today:
 - Can give deep **insight** to real physics of the sound propagation in ducts, 3D data available
 - Simulation of complex geometries and with complex flow conditions
 - **Cost effective** parameter studies
 - **GPUs** will allow optimization in near future
 - Validation and analysis of each specific result remains essential part of the task.
- Experiments cannot be avoided:
 - Accurate experiments are required more than ever to provide **model parameters** and input for CAA.

Thank You!

The DRP based CAA solver TUBA

- 3D finite difference based time-domain method
 - Spatial discretization:
Dispersion-Relation-Preserving Scheme (DRP), [TAM & WEBB, 1993]
 - 7-point central difference scheme, 4th order optimized
- Time discretization:
 - 5/6, resp. 5-stage Runge-Kutta scheme, [HU, 1996 / CARPENTER, ET AL, 1994]
 - Storage optimized (2N-storage), [STANESCU & HABASHI, 1998]
- Spatial low-pass Filter
 - Up to 10th order, 11 point filter stencil
 - Secure suppression of no-resolved short waves
 - Low dissipation of waves resolved above 7 PPW
- TUBA – One method several mathematical models:
 - 2D, modal axisymmetric (AXI), fully 3D
 - Linearized Euler equations (LEE), nonlinear Euler equations (PENNE), isentropic/non-isentrop variants
 - Parallel Implementation (MPI)

➤ Boundary Conditions:

- Source Boundary
 - Analytical source model (modes, frequencies, amplitudes)
 - Modal analyses of experimental data
 - Modal matching approach of CFD data [OVENDEN & RIENSTRA, 2004]
 - Direct input of unsteady CFD data
- Non-reflecting boundary conditions
 - Radiation/outflow condition [BAILLY & BOGEY, 2002]
 - Sponge Layer [ISRAELI & ORZAG, 1981]
 - Perfectly Matched Layer (PML) [HU, ET AL. 1996/2001]
- Wall boundary conditions
 - Hard wall
 - Lined wall, Extended Helmholtz Resonator (EHR) [RIENSTRA, 2006]
- Structured multi domain method
- Overset Grid Method (Chimera)

Anwendung des EHR Modells

Das erweiterte Helmholtz-Reosonator (EHR)- Impedanzmodell im Zeitbereich

- EHR-Modell (Rienstra 2006)
 - Fünf Parameter-Modell für den Frequenzgang der Impedanz
 - SDOF (Totzeit T_l)
 - Massenreaktanz der Deckschicht m_f und Reaktanz des Volumens, sowie offene Fläche in β berücksichtigt
 - Dämpfung der Deckschicht R_f und des Volumens ε berücksichtigt
- Verknüpfung zum klassischen Modell für tiefe Frequenzen über Laurent-Reihe
 - **Verknüpfung zu geometrischen Größen!**

$$Z_{\text{EHR}}(i\omega) = R_f + i\omega m_f - i\beta \cot \left(\frac{1}{2} \omega T_l - i \frac{1}{2} \varepsilon \right)$$

$$= \frac{(R_f + i\omega m_f) (1 - e^{-\alpha}) + \beta (1 + e^{-\alpha})}{1 - e^{-\alpha}},$$

with $\alpha = i\omega T_l + \varepsilon$.

Randbedingung im Zeitbereich:

$$\frac{\partial u'_n}{\partial t}(t) = \frac{1}{m_f} [\mu(t) - (R_f + \beta) u'_n(t)]$$

$$- \underbrace{\frac{1}{m_f} e^{-\varepsilon} [\mu(t - T_l) - (R_f - \beta) u'_n(t - T_l)]}_{\text{storage term}} + e^{-\varepsilon} \frac{\partial u'_n}{\partial t}(t - T_l),$$

Annäherung des Kotangens über eine Laurent-Reihe

$$Z_{\text{EHR}}(i\omega) \approx R_f + \frac{1}{6} \varepsilon \beta + \frac{2\beta\varepsilon}{T_l^2 + (\frac{\varepsilon}{\omega})^2} + i\omega \left(m_f + \frac{\beta T_l}{6} \right) + \frac{1}{i\omega} \left(\frac{2\beta T_l}{T_l^2 + (\frac{\varepsilon}{\omega})^2} \right)$$

- Vorteil der Optimierung im Zeitbereich:
 - Optimierung der Modellparameter für mehrere Frequenzen (z.B. BPF und 2xBPF) liefert direkt die Geometrie des optimalen Liners
 - Aus Erfahrungen mit der *impedance eduction* für Linern im Prüfstand des DLR (AT-TA) kann die Verknüpfung überprüft werden

Verhältnis zwischen korrigierter Hals-Länge und Open Area

$$\frac{L}{\sigma} = m_f + c_0 \frac{1}{6} \beta T_l.$$

Verhältnis aus Fläche und Volumen

$$\frac{S}{V} = \frac{1}{c_0} \frac{2 \beta T_l}{T_l^2 + \frac{\epsilon^2}{\omega^2}}.$$

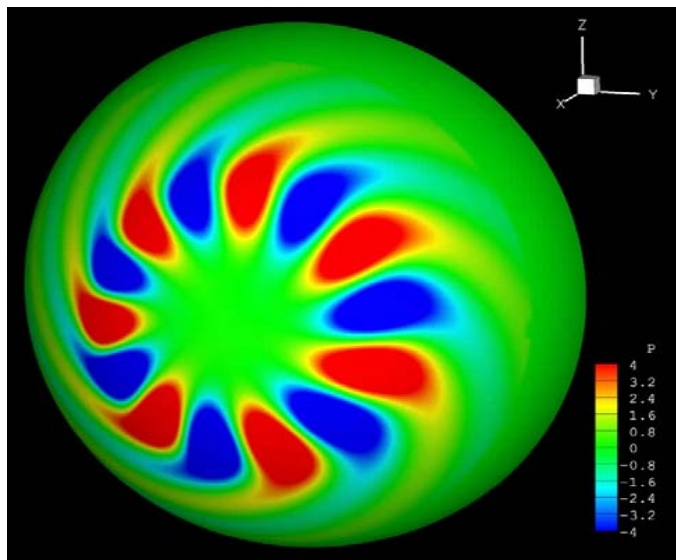
Näherung für die Dämpfung ($\text{Re}\{Z\}$)

$$R_f + \frac{2 \beta \epsilon}{T_l^2 + \left(\frac{\epsilon}{\omega}\right)^2} + \frac{1}{6} \beta \epsilon$$

Beispiel für eine 3D-Rechnung mit Liner

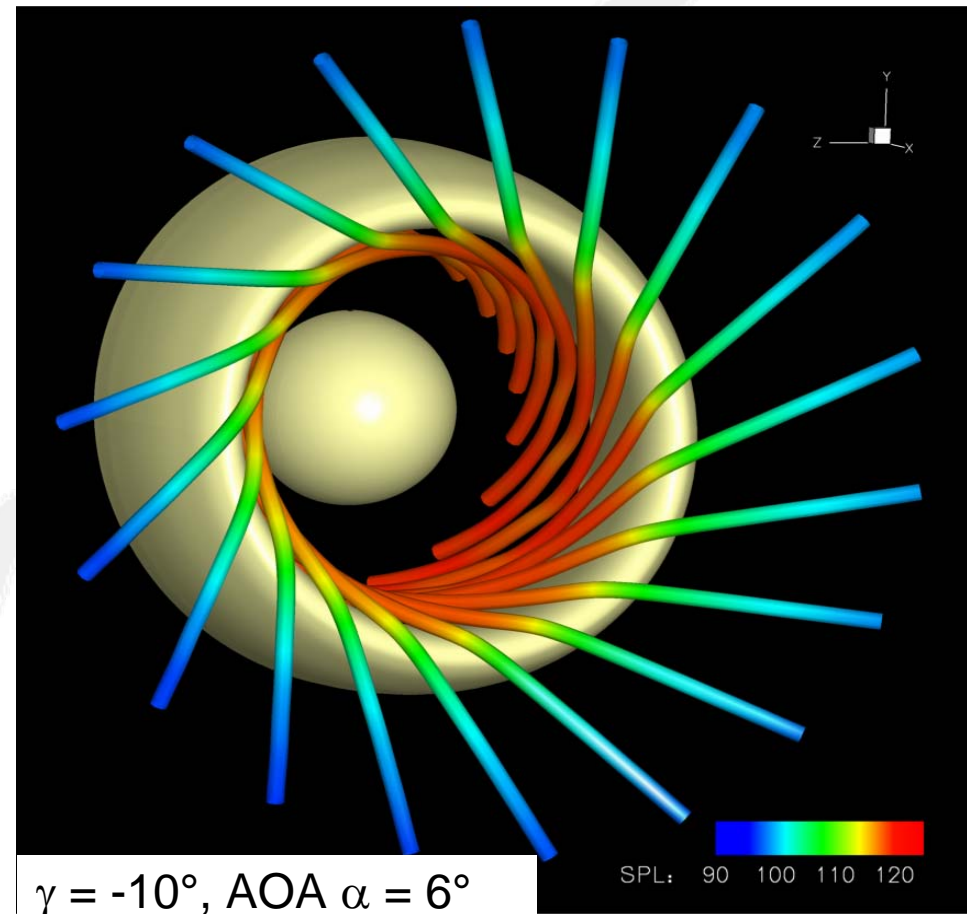
3D scarfed Intake: Auswertung der Schallintensität

- Stromlinien der akustischen Intensität zeigen den Weg des Schalls durch den Einlauf auf.
- ▶ Beeinflussung der Abstrahlung nach unten muss dem Pfad des nach unten abgestrahlten Schalls folgen.
- ▶ (7,0) mode, $f = 1$ kHz
- ▶ $M_{\text{flight}} = 0.2$ $M_{\text{fan-plane}} = 0.35$



**Momentan-
aufnahme des
Schalldrucks**

FW/H
Austauschfläche,
Abstand 1m



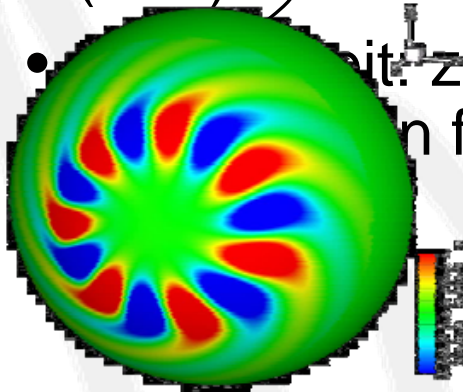
Anwendung der akustischen Analogie von Ffowcs-Williams & Hawkings

➤ Akustische Analogie nach Ffowcs-Williams & Hawkings (FW-H)

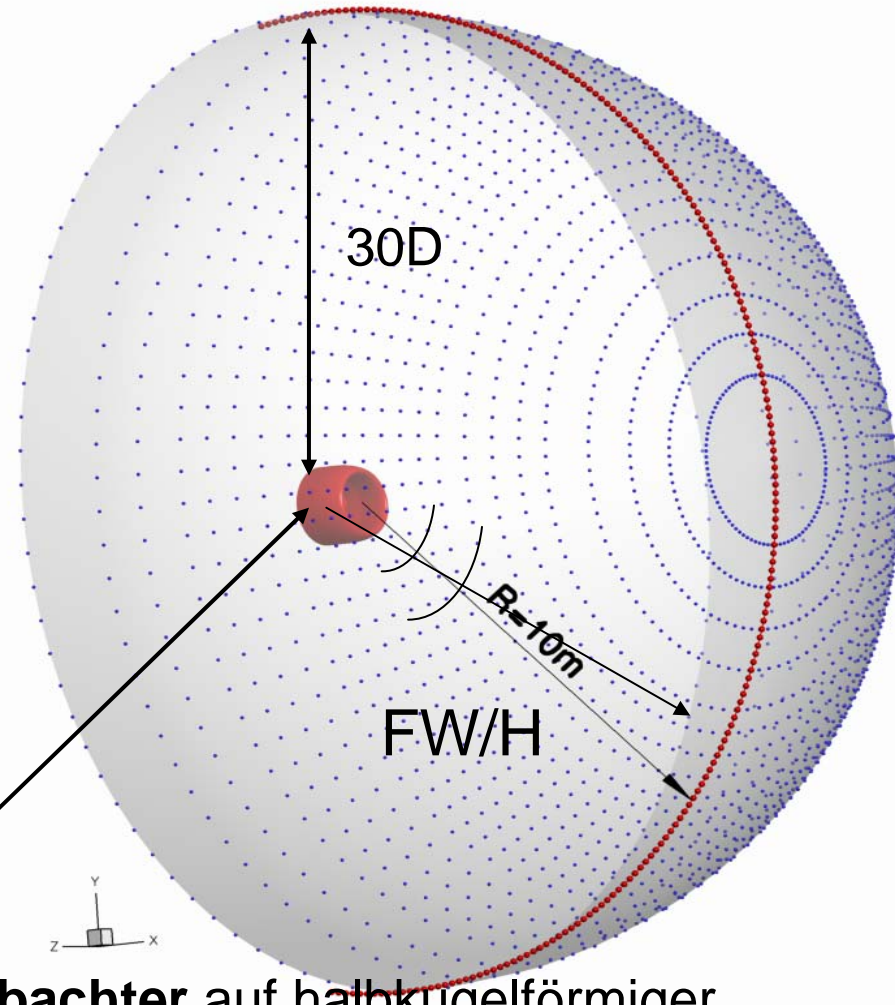
- Ausnutzung der Zeit-Periodizität möglich (nicht notwendigerweise harmonisch!)

- Parallele Implementierung Austauschfläche (MPI)

- Rechenzeit: z.B. 1h auf 4 Prozessoren für 10^4



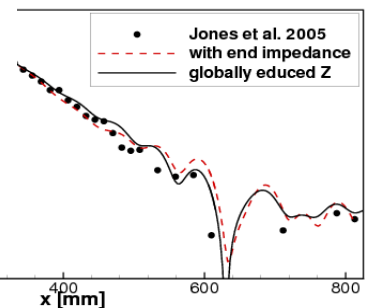
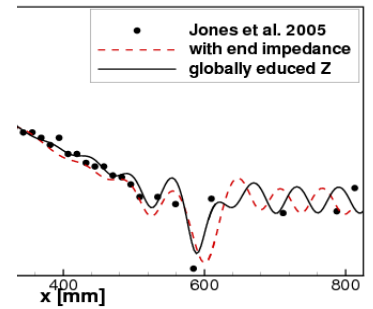
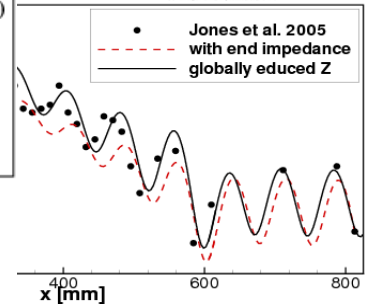
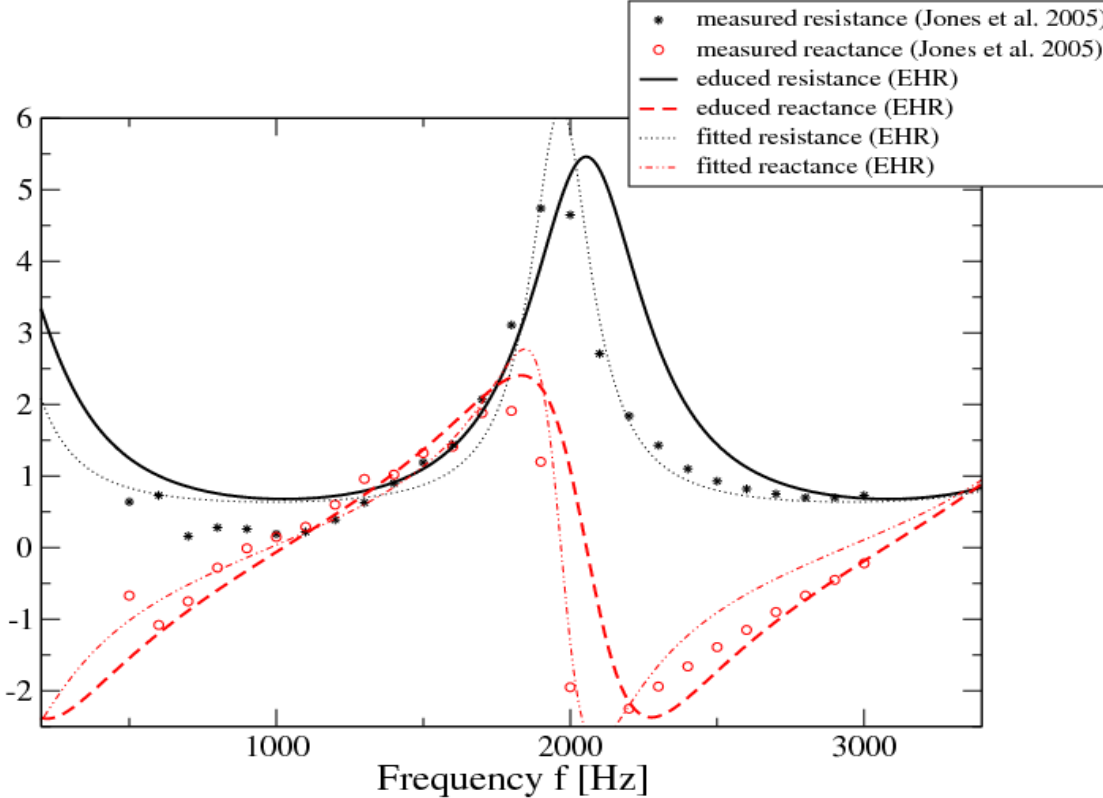
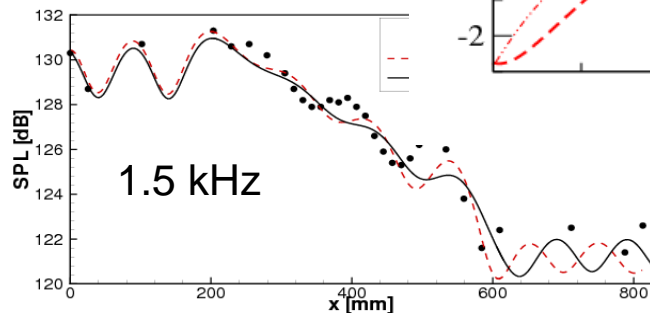
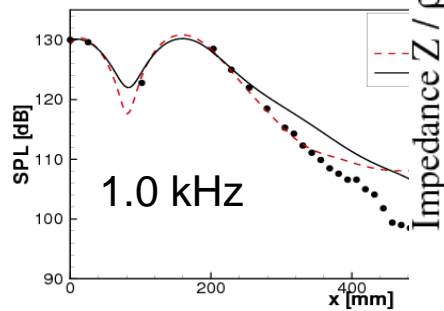
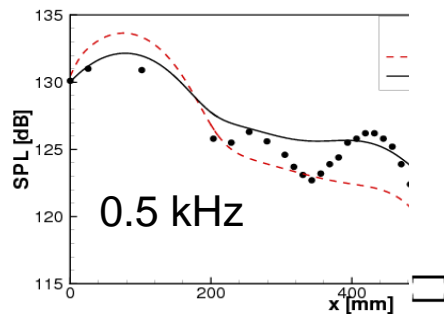
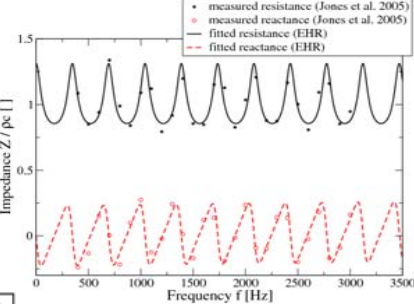
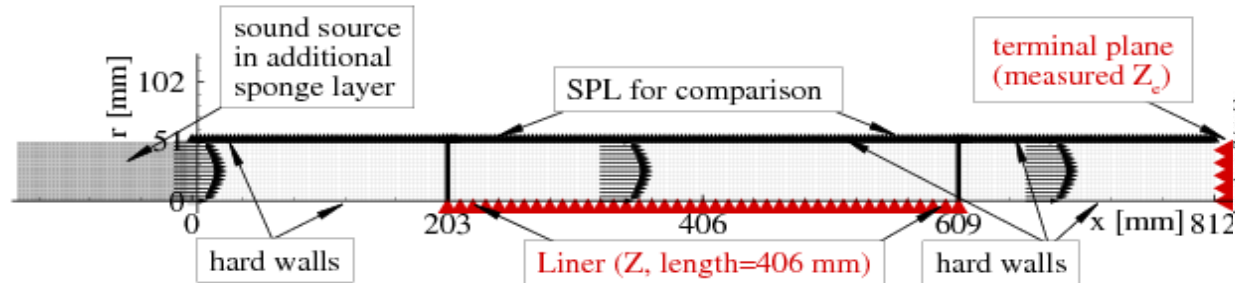
3D CAA



Beobachter auf halbkugelförmiger Fläche vor dem Einlauf

Validierung des Verfahrens mit Liner

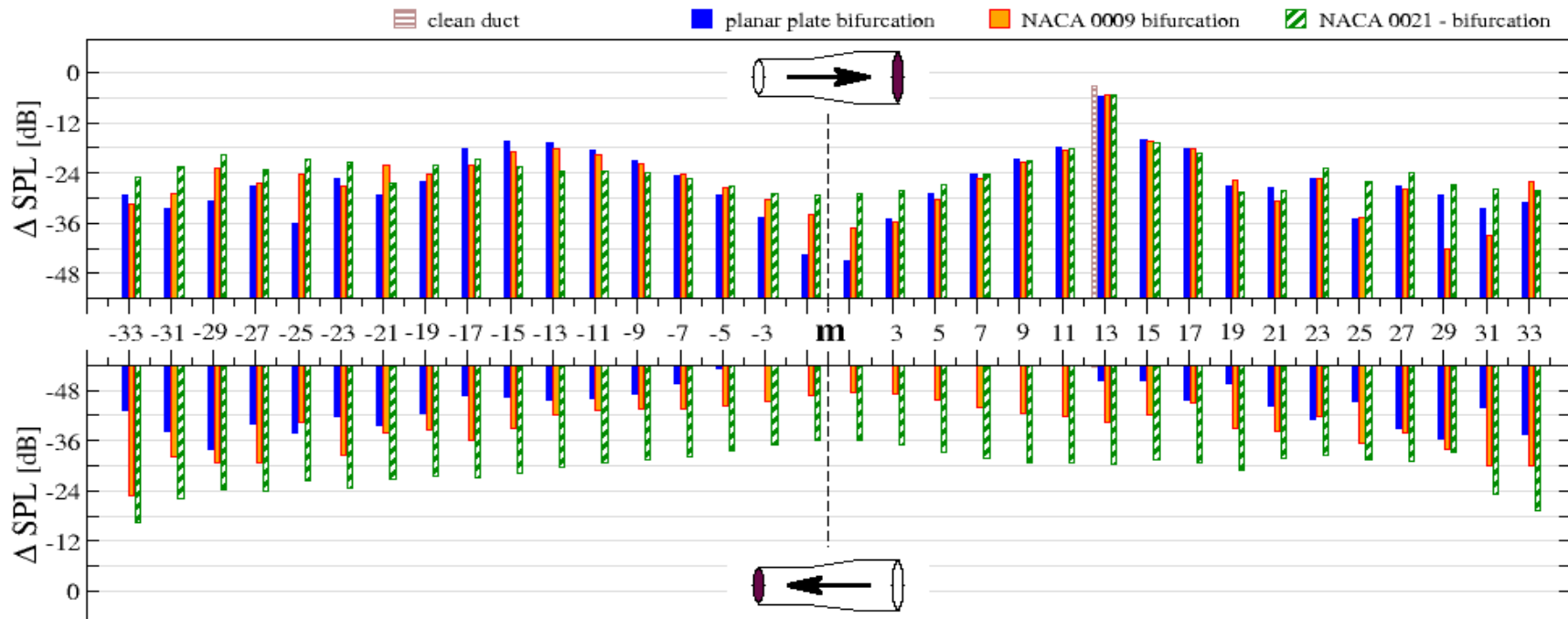
NASA GIT Experiment, AIAA Journal 2007



➤ Folgende Folien sind aus dem Original entfernt

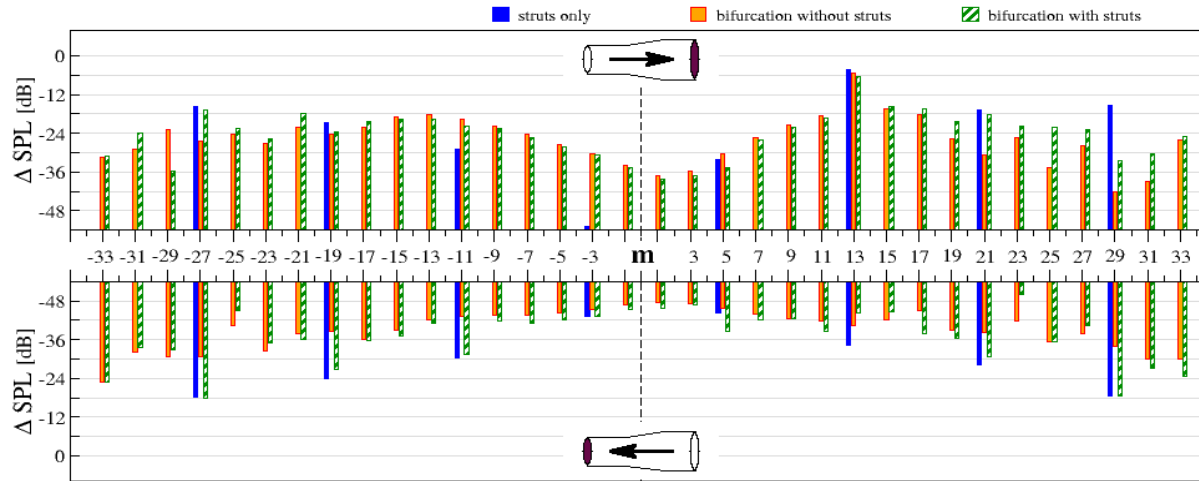
Modal analysis of installation effects – Effect of the bifurcation profile width

- Correspondingly to the theory of TYLER & SOFRIN (1962) the two splitter cause a scattering of the source mode into every second propagation capable mode ($m \pm 2S$; S is a natural number)
- Source mode (13,0) keeps dominant
- Scattering cause a reduction of the source mode amplitude by $\sim 2.5\text{dB}$
- Scattered modes are at least 12dB quieter than source mode

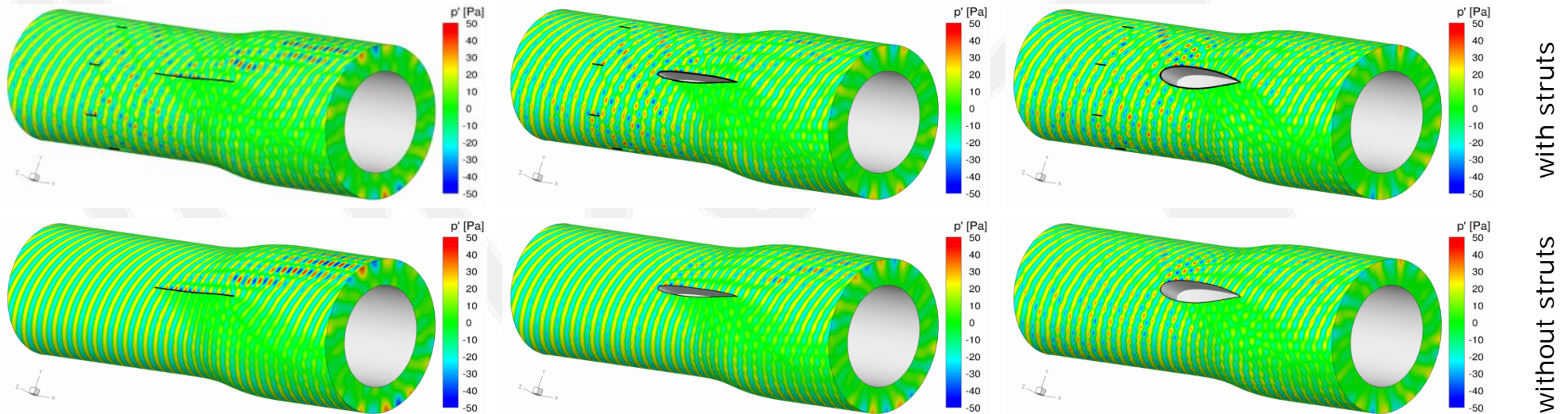


Mode analysis of installation effects – Combination of struts and splitter

► Combination of Splitters and struts



- Similar mode spectrum for cases with and without struts
- Only modes excited in both configurations change noticeable, especially near the inlet



Validation of each specific CAA result

- Finite difference based CAA methods have no conservation properties.
 - It is easy to get 20dB of numerical dissipation by choosing a wrong setting for the filter or having under resolved meshes in crucial regions.
- How to monitor the correctness of a CAA result and secure the interpretations and conclusions drawn from the numerical experiment?
- Suggested solution: Assessment of the acoustic energy conservation.
- Problem: perturbation equations are applied as CAA-model, for which no conservation properties can be stated in general:
 - Conservation of energy, momentum and mass applies only for the whole flow field.
 - Most general formulation is perturbation energy conservation for the acoustic energy by Morfey [1971] applies in a potential flow field with acoustic perturbation modes only.
 - However, allows limited validation of the results.

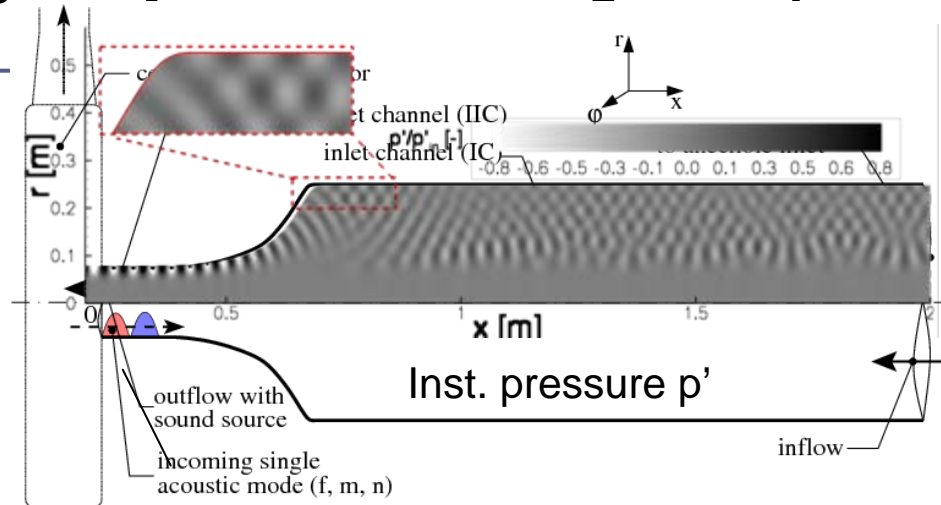
- A *by view* analysis of the solution quality by watching the sound pressure cannot capture scattered modes and transient processes as well.
- Acoustic intensity is calculated accordantly to the definition of [MORFEY, 1971] during the simulation an averaged in time
(e.g. over one period of a harmonic oscillation)
- Entire energy flux at arbitrary axial positions is given as integral of the time averaged acoustic intensity over the cross-section:

$$P_a(x) = \int_{A(x)} \underline{I}_a \cdot d\underline{A}$$

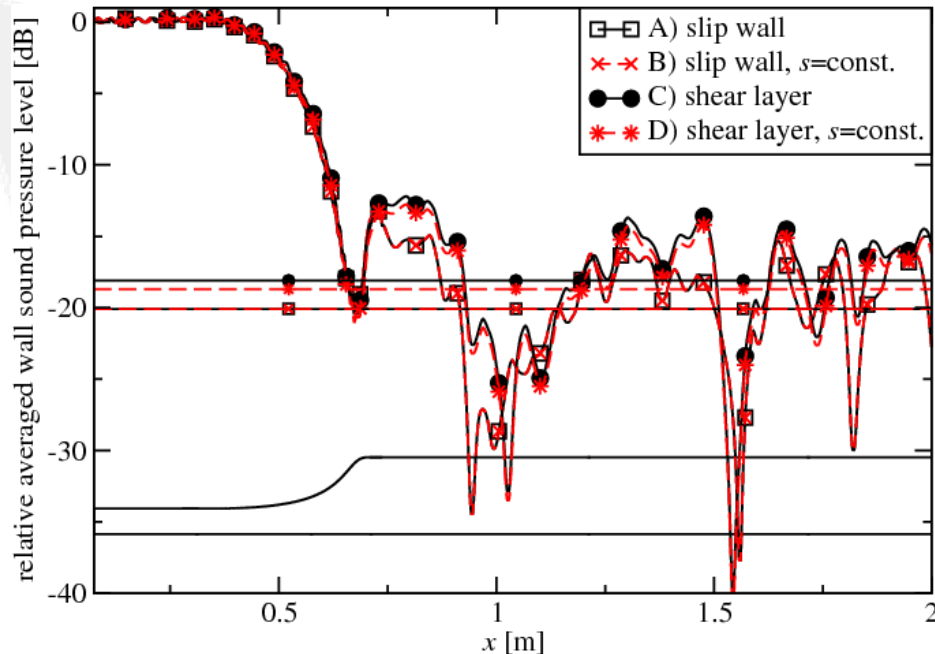
- Acoustic power at all cross sections should be constant in source and sink free hard walled ducts (with potential flow and acoustic modes only).
- With liner, there is a power flux into the lined wall
 - Global conservation of the incoming and outgoing acoustic energy should apply for each control volume.
- Acoustic power is suitable to quantify numerical dissipation.
- Divergence of the time averaged acoustic intensity is a measure for the local source strength → identifies regions with large dissipation.

Intensity analysis for a centrifugal compressor inlet duct

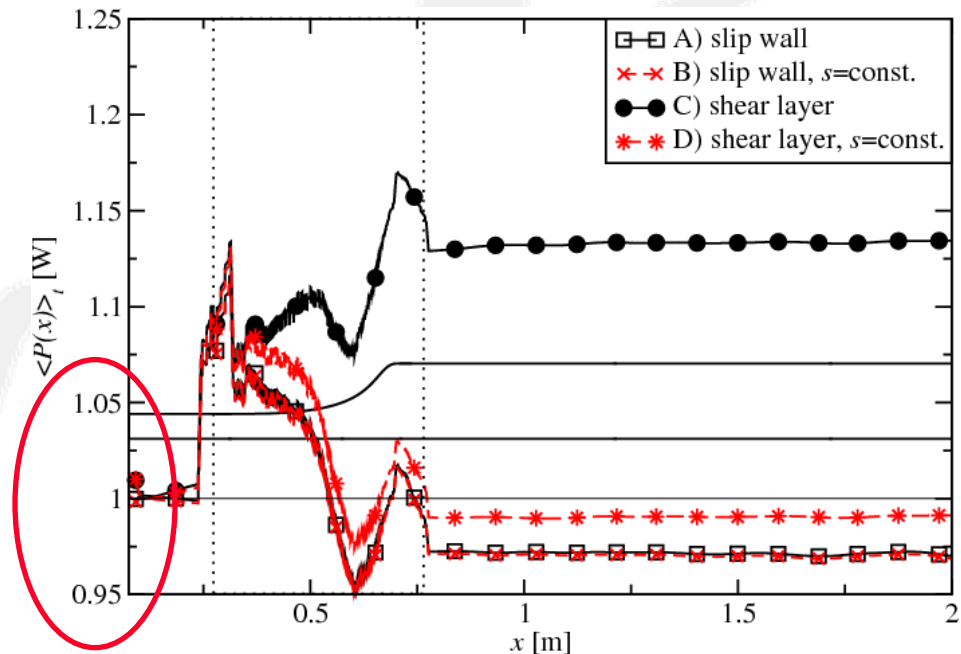
Excitation of a single $m=13, n=0$ mode and strong mode scattering at inlet nozzle



Background flow: realistic boundary layer (Stokes) and potential flow (Euler) at $Ma=0.3$ at the inlet



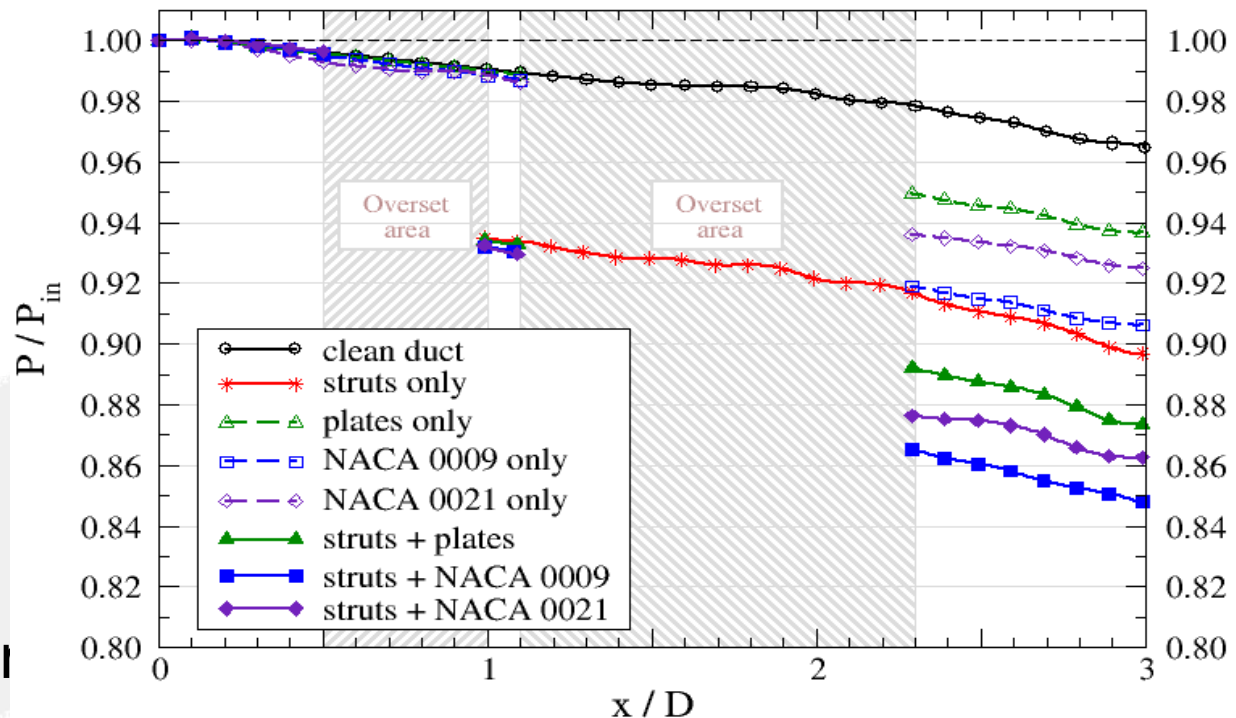
Wall sound pressure level shows a „loss“ of ~20 dB! – Would you expect this result can be correct?



Acoustic power is dissipated or amplified by up to 15% (3% for potential flow) → errors are small.

Intensity analysis of the solution quality in a bifurcated S-duct with Chimera applied

- Back to the S-duct configuration
- Acoustical power distribution in axial direction
 - 3% general dissipation by discretization errors
 - 3-12% extra dissipation by interpolation errors and stronger grid distortion
- Overall error:
 - up to 15% loss of acoustical power
 - up to 0.7dB loss of sound pressure
- Acceptable for industrial application

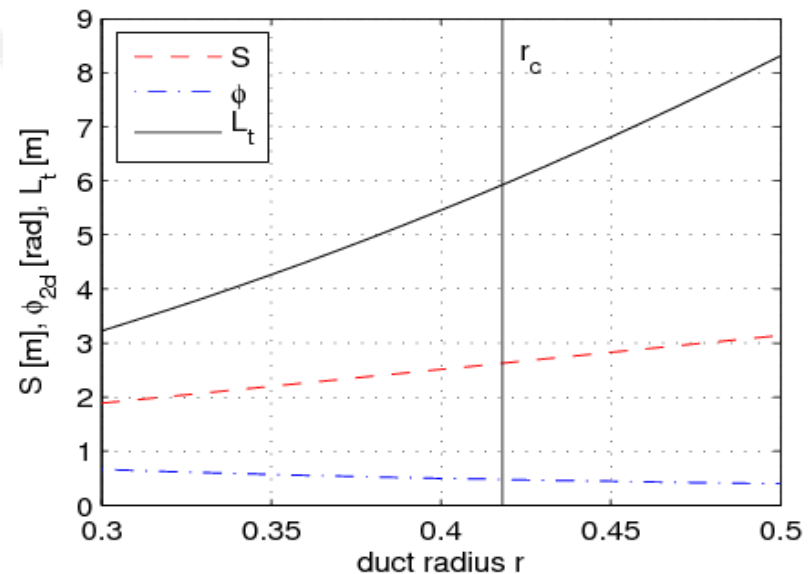
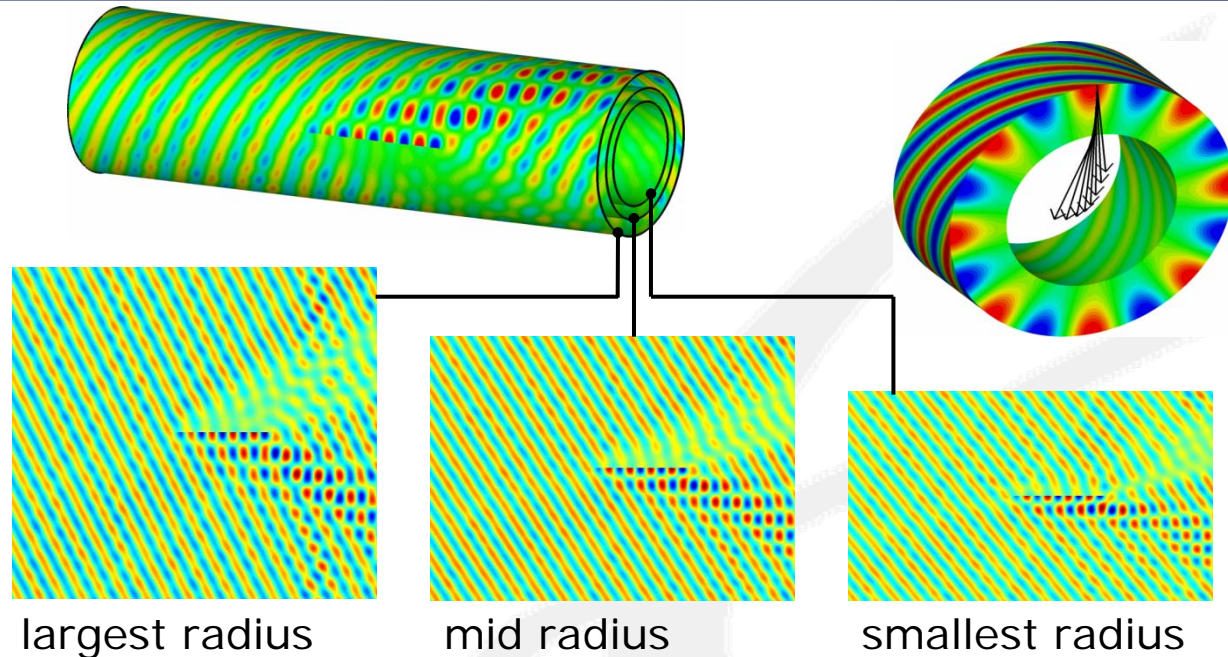


Development of an analytical model from the CAA result – unroll radius

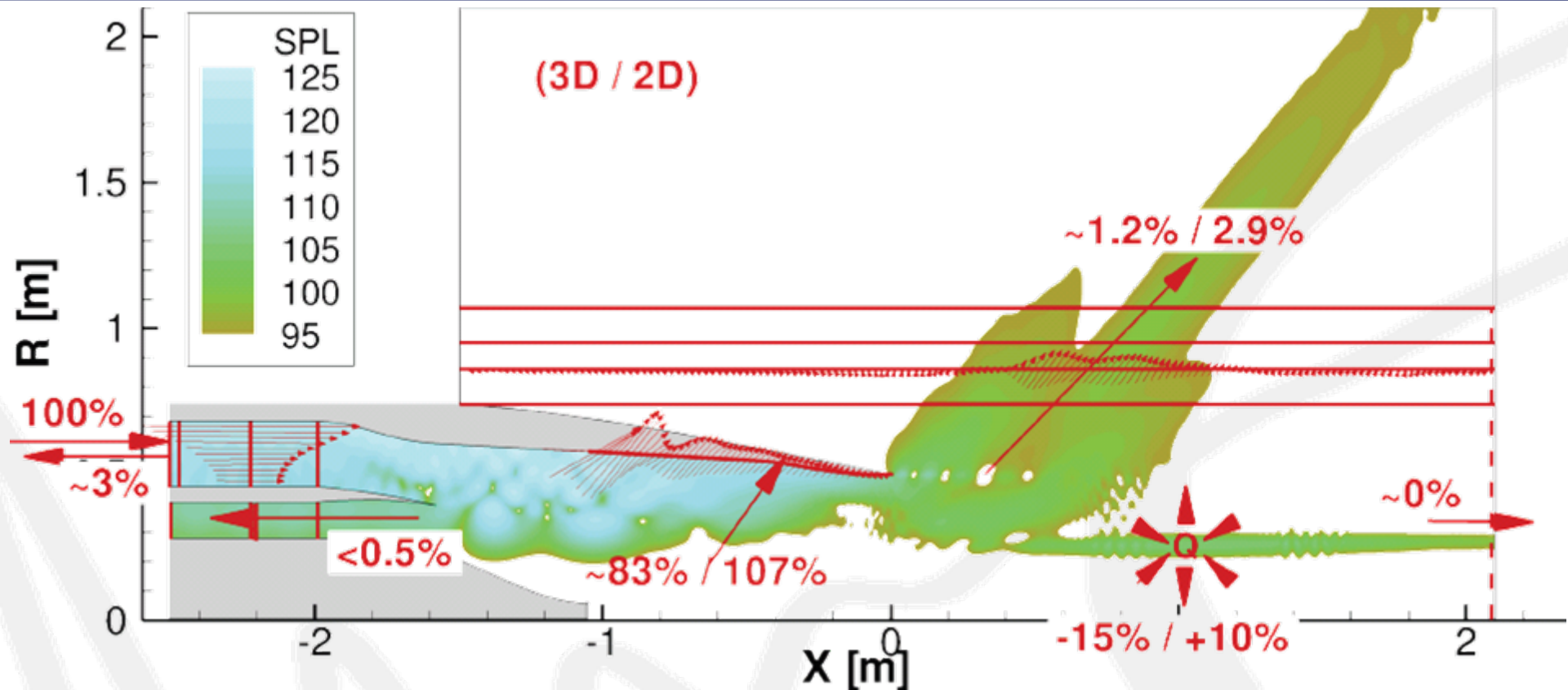
At which radius should the duct be unrolled?

- Reversion length depends on
 - Propagation angle ϕ
 - Splitter distance S
- Caustic radius [Chapman, 1994] and outer radius are limits of the perfect radius to unroll the duct
- Unroll radius depends on the geometry and the mode (eigenvalue)
- Caustic radius:

$$r_c = \frac{|m|}{\sigma_{mn}} R$$



Radiation from a coaxial nozzle (2D/3D) – Detailed analysis of the acoustic energy flux

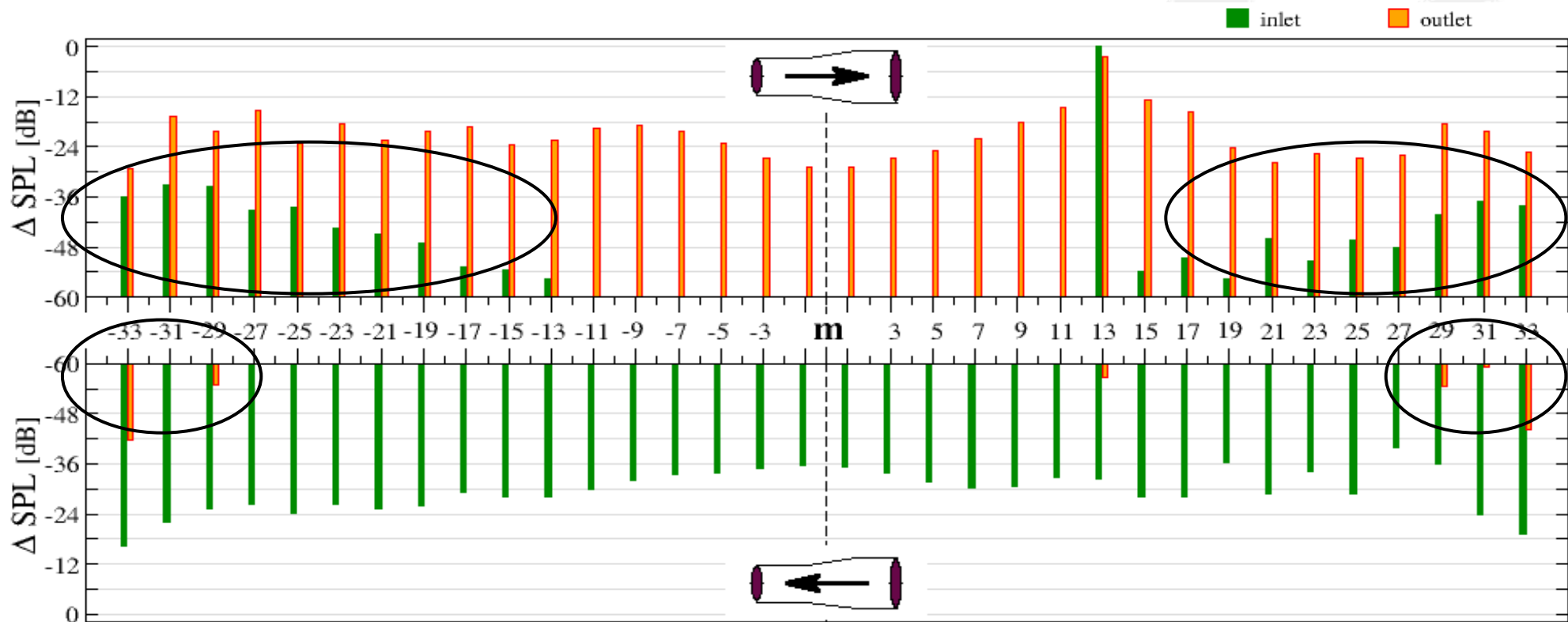


- Global energy balance:
 - Source input is 100 %
 - Outgoing fluxes over bypass duct, core duct, liner and nozzle exit
 - Production and annihilation of ac. energy: non-isentropic, non-potential flow and numerical dissipation
- In the hard walled case all acoustic energy (104 %) is radiated.

The potential of GPU computing

Modal analysis provides validation of the non-reflective properties of the BC's

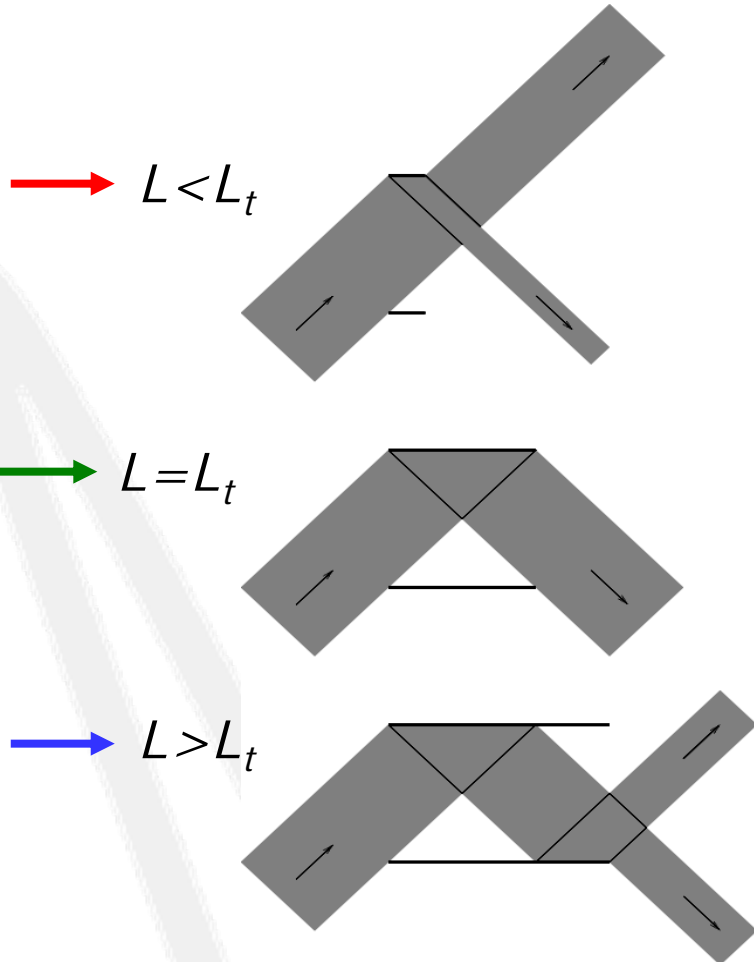
- Visible non-physical reflections in a clean duct:
 - At inlet and outlet boundary, for high azimuthal mode orders
 - Amplitudes more than 12dB lower than physical modes and 30dB lower than source mode → small impact on the entire solution
 - **PML** at outlet performs **better** then sponge layer at inlet



How to describe the sound field in the bifurcation?

Derivation of a theoretical model from the CAA results – Definition of a mode reversion length

Definition of a characteristic length L_t , for which source mode is reversed totally ($m \rightarrow -m$)!



$$L_t = \frac{S}{\tan^{-1} \phi_{2d}}$$

$$f(L) = \begin{cases} 1 - \frac{L}{L_t} & \text{for } 0 \leq L \leq L_t \\ -1 + \frac{L}{L_t} & \text{for } L_t < L \leq 2L_t \end{cases}$$

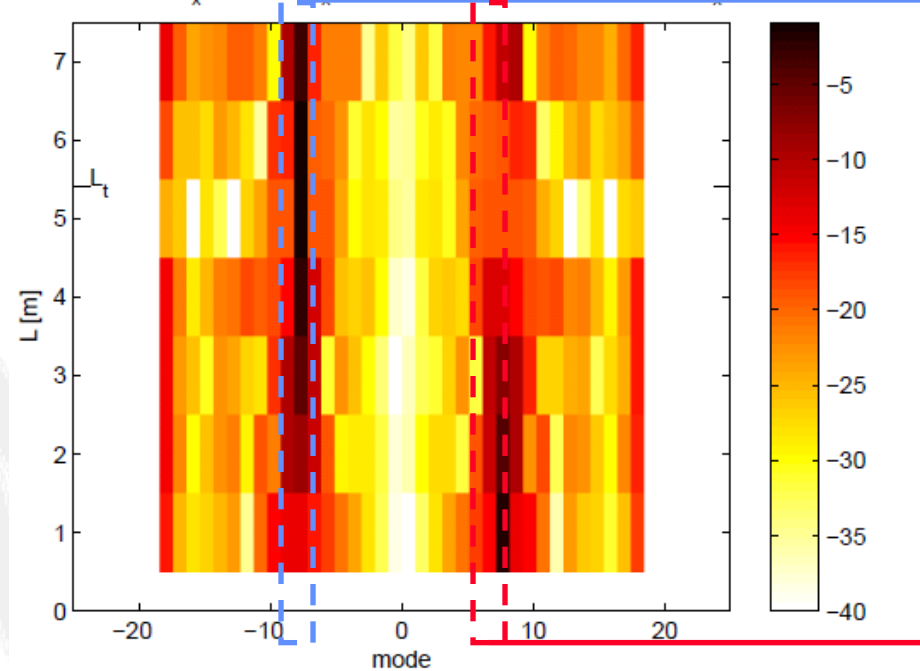
$$g(L) = 1 - f(L)$$

$f(l)$ ist the amplification factor for the initial source mode (m),

$g(l)$ denotes the amplification of the reversed mode ($-m$)

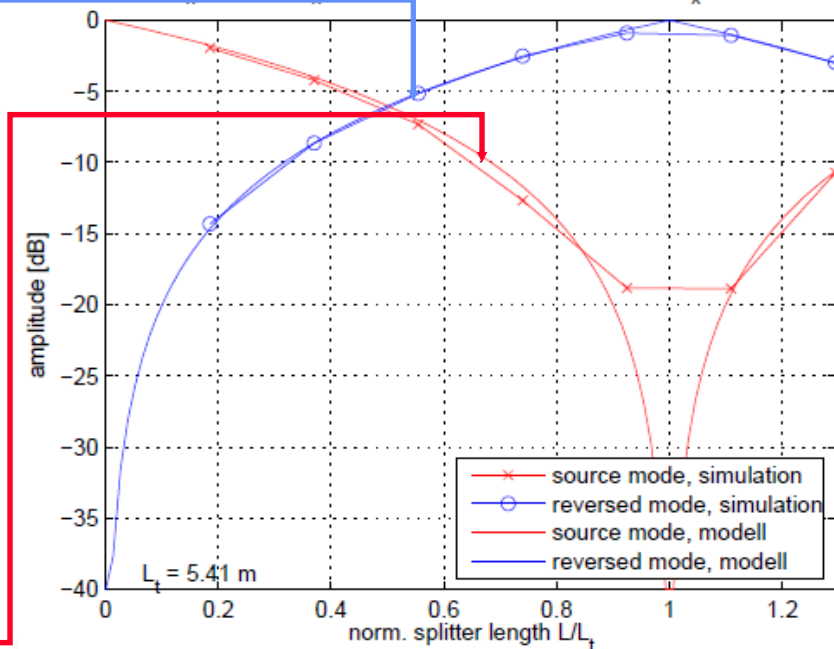
Attenuation of modes by hard infinite thin splitter – analytical model based on CAA simulations

amplitude [dB], $\phi_x = 28.1^\circ$, $\psi_x = 28.1^\circ$, $mn = (8,1)$, $f = 2200$ Hz, $M_x = 0.00$



Mode amplitude after splitter over splitter length and mode

$\phi_x = 28.1^\circ$, $\psi_x = 28.1^\circ$, $mn = (8,1)$, $f = 2200$ Hz, $M_x = 0.00$



Mode amplitude for $|m|=8$ behind splitter over splitter length

- Acoustic energy oscillates periodically between positive and negative modes with increasing splitter length.
- Periodicity of spectrum over splitter length